Heavy Flavour Cross Sections Measurements

$$\left.\begin{array}{lcl} m_u & \approx & 3 \quad MeV \\ m_d & \approx & 5 \quad MeV \\ m_s & \approx & 100 \quad MeV \end{array}\right\} \text{very light } m<\Lambda_{QCD}$$

$$\left.\begin{array}{lcl} m_c & \approx & 1300 \quad MeV \\ m_b & \approx & 4200 \quad MeV \end{array}\right\} \text{Flavours considered}$$

$$m_t \approx 170000 \quad MeV \longrightarrow \text{very heavy}$$

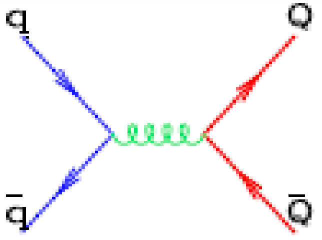
Heavy Flavour Cross Sections at LO

The leading-order process for the production of heavy quark Q of mass m in hadron collisions:

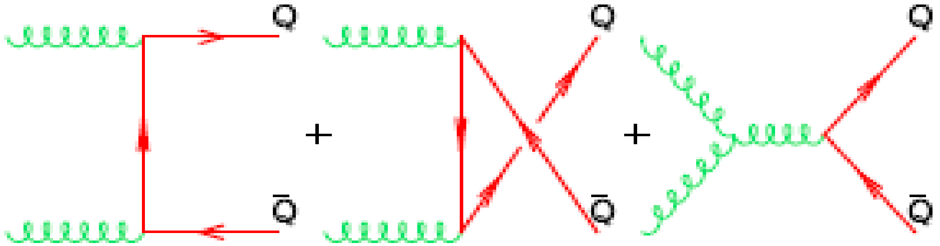
$$\begin{array}{ll} (a) & q(p_1) + \overline{q}(p_2) \rightarrow Q(p_3) + \overline{Q}(p_4) \\ (b) & g(p_1) + g(p_2) \rightarrow Q(p_3) + \overline{Q}(p_4) \end{array}$$

Where the four momenta of the partons are given in brackets.

The Feynman diagrams are:



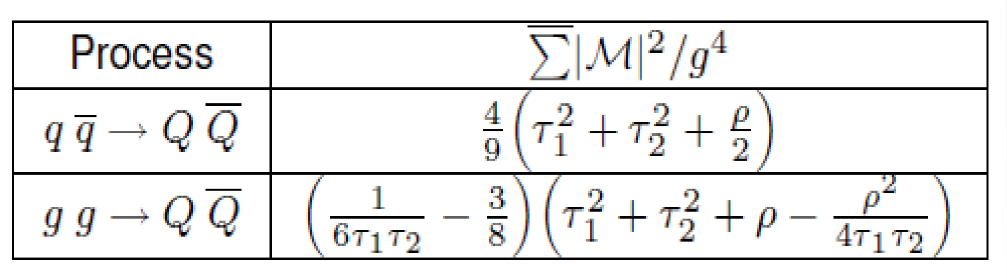
(a)



(b)

Heavy Flavour Cross Sections at LO

The invariant matrix elements squared averaged over initial and final color and spin



Process	$\overline{\sum}\lvert\mathcal{M}\rvert^2/g^4$
$q\,\overline{q} \to Q\,\overline{Q}$	$\frac{4}{9}\left(\tau_1^2+\tau_2^2+\frac{\rho}{2}\right)$
$g\,g \to Q\,\overline{Q}$	$\left(\frac{1}{6\tau_1\tau_2}-\frac{3}{8}\right)\left(\tau_1^2+\tau_2^2+\rho-\frac{\rho^2}{4\tau_1\tau_2}\right)$

Where it has been introduced the notation:

$$\tau_1 = \frac{2p_1.p_3}{\hat{s}},\quad \tau_2 = \frac{2p_2.p_3}{\hat{s}},\quad \rho = \frac{4m^2}{\hat{s}},\quad \hat{s} = (p_1+p_2)^2$$

The short-distance cross section is obtained from the invariant matrix element:

$$d\hat{\sigma}_{ij} = \frac{1}{2\hat{s}}\frac{d^3p_3}{(2\pi)^3 2E_3}\frac{d^3p_4}{(2\pi)^3 2E_4}(2\pi)^4\delta^4(p_1+p_2-p_3-p_4)\overline{\sum}\lvert\mathcal{M}_{ij}\rvert^2.$$

Heavy Flavour Cross Sections at LO

In terms of rapidity $y = \frac{1}{2}\ln((E+p_z)/(E-p_z))$ and of transverse momentum P_T the relativistically invariant space volume element of the final state heavy quark is :

$$\frac{d^3p}{E} = dy\, d^2p_T$$

The invariant cross section may be written at LO:

$$\frac{d\sigma}{dy_3 dy_4 d^2p_T} = \frac{1}{16\pi^2\hat{s}^2}\sum_{ij} x_1 f_i(x_1,\mu^2)\, x_2 f_j(x_2,\mu^2)\overline{\sum}|M_{ij}|^2$$

x_1 and x_2 are fixed if transverse momenta and rapidity of the outgoing heavy quarks are known. In the CM of the incoming hadrons we can write

$$\begin{aligned}
p_1 &= \frac{1}{2}\sqrt{s}(x_1,0,0,x_1)\\
p_2 &= \frac{1}{2}\sqrt{s}(x_2,0,0,-x_2)\\
p_3 &= (m_T\cosh y_3, p_T, 0, m_T\sinh y_3)\\
p_4 &= (m_T\cosh y_4, -p_T, 0, m_T\sinh y_4)
\end{aligned}$$

Heavy Flavour Cross Sections at LO

Applying the energy and momentum conservation

$$\begin{aligned} x_1 &= \frac{m_T}{\sqrt{s}}\left(e^{y_3}+e^{y_4}\right) \\ x_2 &= \frac{m_T}{\sqrt{s}}\left(e^{-y_3}+e^{-y_4}\right) \\ \hat{s} &= 2m_T^2(1+\cosh\Delta y). \end{aligned}$$

$m_T = \sqrt{m^2 + p_T^2}$ transverse mass of the heavy quarks

$\Delta y = y_3 - y_4$ rapidity difference between heavy quark

With this notation the matrix elements

$$\overline{\sum}|\mathcal{M}_{q\bar{q}}|^2 = \frac{4g^4}{9}\left(\frac{1}{1+\cosh(\Delta y)}\right)\left(\cosh(\Delta y)+\frac{m^2}{m_T^2}\right), \quad \sim \text{costant}$$

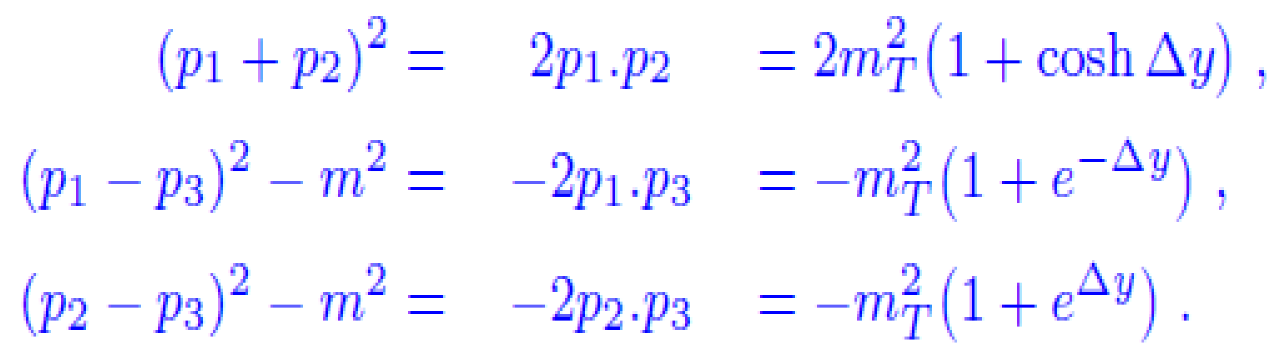
$$\overline{\sum}|\mathcal{M}_{gg}|^2 = \frac{g^4}{24}\left(\frac{8\cosh(\Delta y)-1}{1+\cosh(\Delta y)}\right)\left(\cosh(\Delta y)+2\frac{m^2}{m_T^2}-2\frac{m^4}{m_T^4}\right) \quad \sim \exp(-\Delta y)$$

Low contribution at high Δy and dominant contribution for Δy<1

Heavy quarks produced by light quark are more correlated in rapidity respect to those produced by gluon-gluon fusion.

Applicability of Perturbation Theory

The propagators in the diagrams:



$$
\begin{aligned}
(p_1+p_2)^2 &= \quad 2p_1.p_2 &&= 2m_T^2\left(1+\cosh\Delta y\right) \,, \\
(p_1-p_3)^2-m^2 &= -2p_1.p_3 &&= -m_T^2\left(1+e^{-\Delta y}\right) \,, \\
(p_2-p_3)^2-m^2 &= -2p_2.p_3 &&= -m_T^2\left(1+e^{\Delta y}\right) \,.
\end{aligned}
$$

Are off-shell by a quantity of the order of m^2 so the perturbation theory should be applicable.

This is valid until the mass m is larger of Λ_{QCD}.

The question is if the bottom and charm mass are large enough.

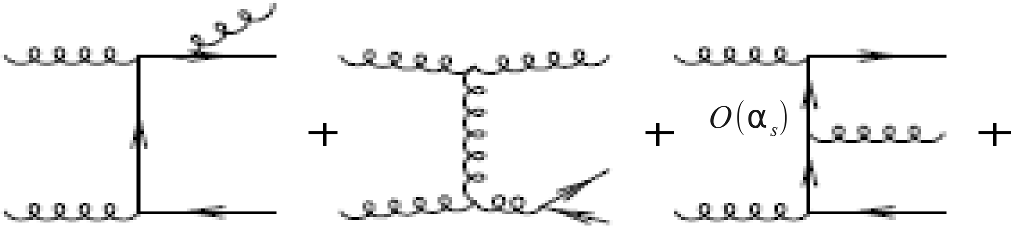
Heavy Flavour Cross Sections at higher-order

At NLO, $O(\alpha_s^3)$, the production cross section of the heavy flavor quark of mass m:

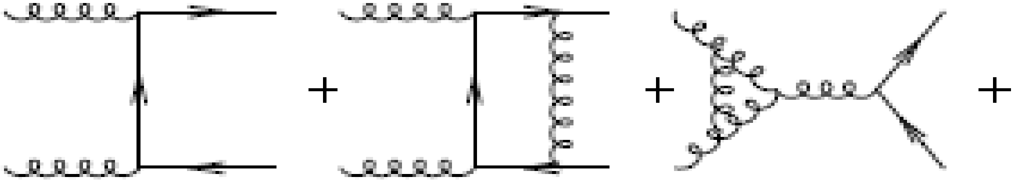
$$\sigma(S)=\sum_{i,j}\int dx_1 dx_2\, \hat{\sigma}_{ij}(x_1x_2S,m^2,\mu^2)F_i(x_1,\mu^2)F_j(x_2,\mu^2)$$

Where $\hat{\sigma}_{i,j}(\hat{s},m^2,\mu^2)=\sigma_0 c_{ij}(\hat{\rho},\mu^2)$ $\hat{\rho}=4m^2/\hat{s},\bar{\mu}^2=\mu^2/m^2,\sigma_0=\alpha_S^2(\mu^2)/m^2$

Examples of higher order diagrams



Real emission diagrams



Virtual emission diagrams

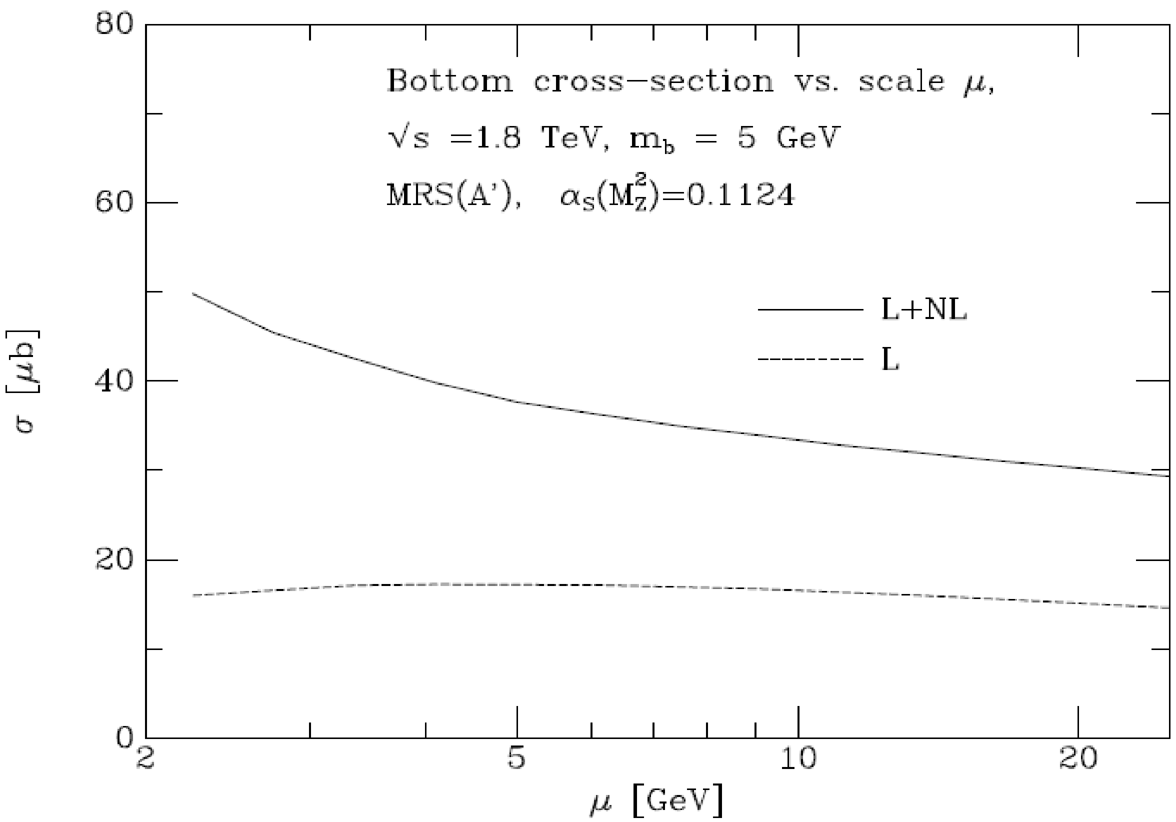
Heavy Flavour Cross Sections at higher-order

There are several dependences.

1) Scale, μ:
  - PDF following the DGLAP equations
  - running coupling constant
  - short-distance cross section: if we perform a calculation to $O(\alpha_s^3)$, the variation of the scale contributes: $\mu^2 \frac{d}{d\mu^2}\sigma = O(\alpha_S^4)$.

The above variations combine in a such a way that the scale dependence is formally small because of higher order in α_s (α_s^4)
This does not guarantee that the numerical value of the cross section is smaller for higher series when varying the scale.

Heavy Flavour Cross Sections at higher-order



LO cross section is almost scale independent because of α_s behavior and increased gluon distribution with μ

NLO is almost 2xLO → large uncertainties on the cross section

Heavy Flavour Cross Sections at higher-order

2) Heavy quark mass m :

$1.2 < m_c < 1.8$ GeV

$4.5 < m_b < 5.0$ GeV

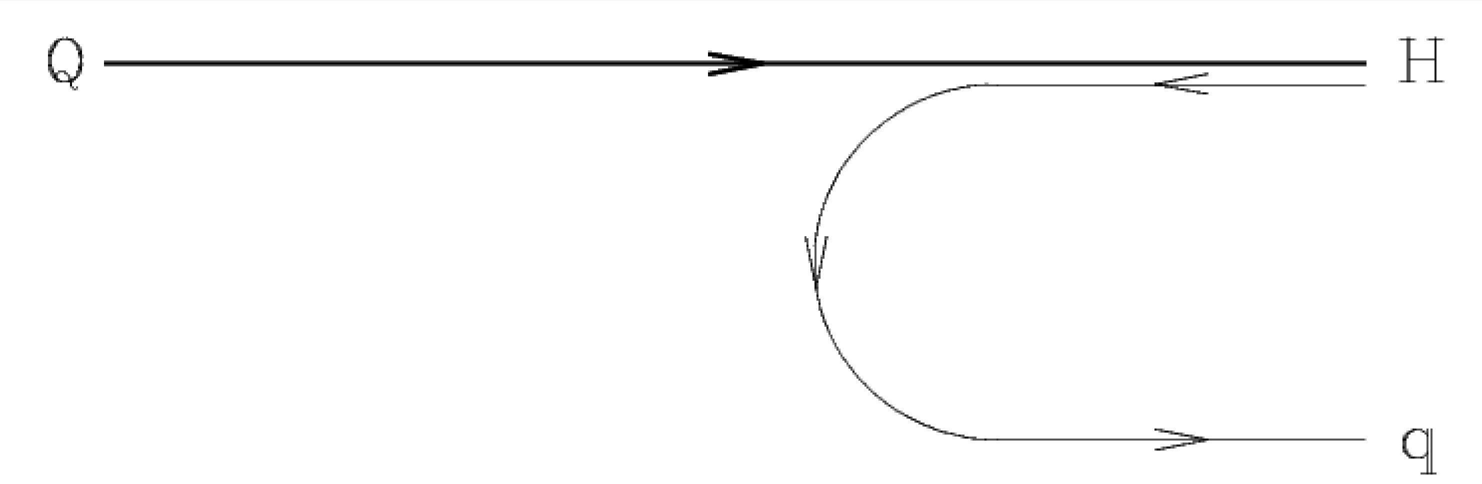
- explicit dependence on $1/m^2$ in the short-distance cross section
- PDF, as m decreases the x value at which the PDF are calculated become smaller and the cross section increases because the parton flux increase
- α_s depends on the the scale μ and Λ

$$\alpha_s(\mu^2)=\frac{1}{b_0 \ln(\frac{\mu^2}{\Lambda^2})}$$

if we take $m/2 < \mu < 2m$ we have problems with the charm because we arrive at $\mu < 1$ GeV where the perturbation theory is not valid → for charm the lower limit for $\mu=2m_c$

Heavy Quarks Fragmentation

Heavy quarks after the production fragment in hadrons.
The model is different of those used for light quarks
the attachment of a light quark to and heavy one Q produce a small deceleration of the heavy quark Q.



Heavy Quarks Fragmentation

An heavy quark Q of momentum P generates a hadron H=Qq of momentum zP. To model this process the energy difference before and after the fragmentation is needed

$$\Delta E = E_Q - E_H - E_q = \sqrt{m_Q^2 + P^2} - \sqrt{m_H^2 + z^2 P^2} - \sqrt{m_q^2 + (1-z)^2 P^2}$$

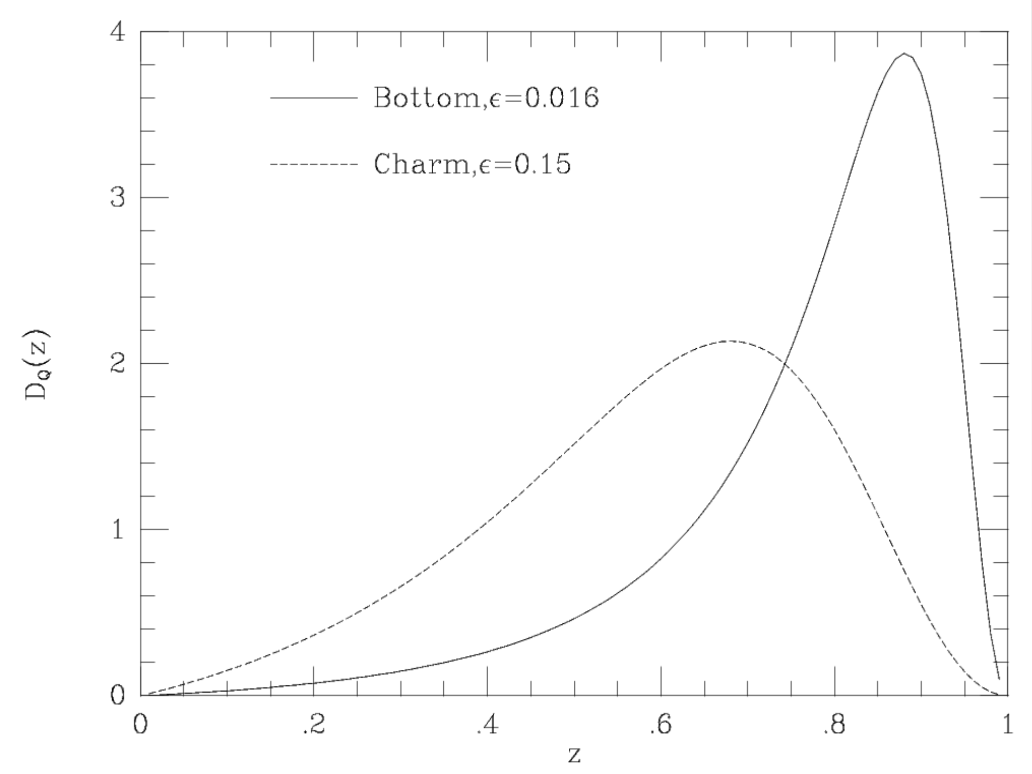
$$\approx \frac{m_Q^2}{2P}\left[1 - \frac{1}{z} - \frac{\epsilon_Q}{1-z}\right] \quad \text{with} \quad \epsilon_Q = \frac{m_q^2}{m_Q^2} \quad \text{and} \quad m_H = m_Q$$

The transition amplitude is $T \sim \frac{1}{\Delta E}$, squaring the amplitude and including a factor 1/z for phase space we obtain the Peterson function for the heavy quark fragmentation

$$D_Q^H(z) = \frac{N_H}{z}\left[1 - \frac{1}{z} - \frac{\epsilon_Q}{1-z}\right]^{-2}$$

N_H is a normalization factor $\sum \int dz\, D_Q^H(z) = 1$

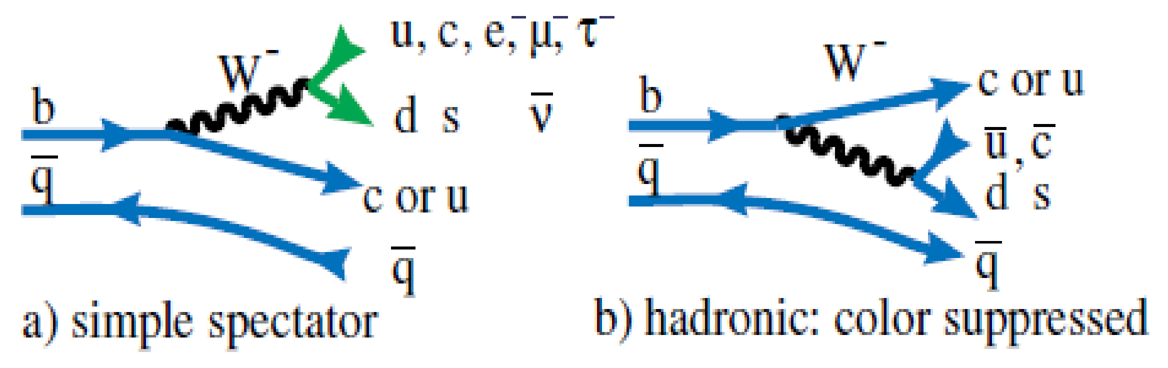
Heavy Quarks Fragmentation



$$\epsilon_Q = \frac{m_q^2}{m_Q^2}$$

is determined by the ratio of quark masses but it is treated as parameter and the best value is obtained from data.

b and c Meson Decay



a) simple spectator

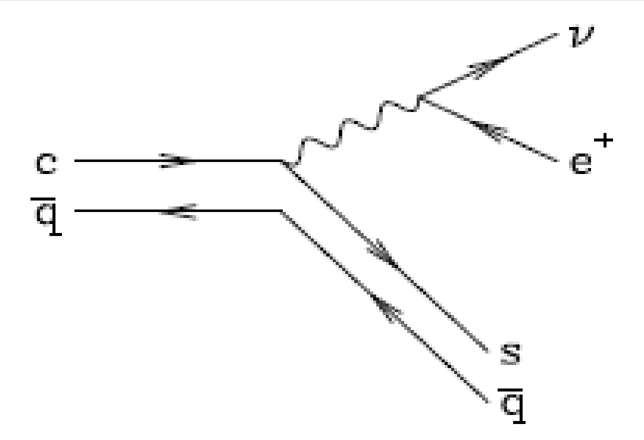
b) hadronic: color suppressed

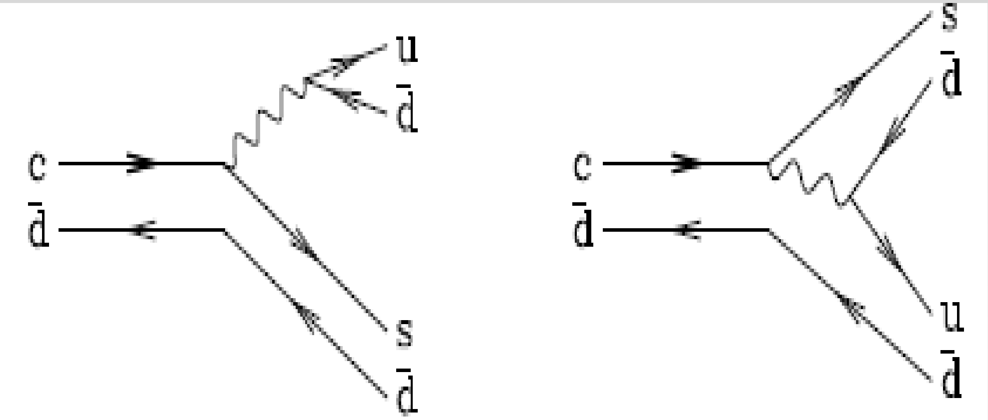
Several different decay modes of the b meson based on the spectator Model where the light quark does not participate.

B can be selected based on:

- lepton (e,μ): B → lνD
- D meson: B → Dπ
- J/ψ: B → J/ψ K

c Meson Decay





Similar diagram governs the charm decay.

The procedure to identify charm meson are:

- lepton (e,μ): D → lνK(π)
- K meson: D → K(n)π

b Meson Cross Sections Measurements

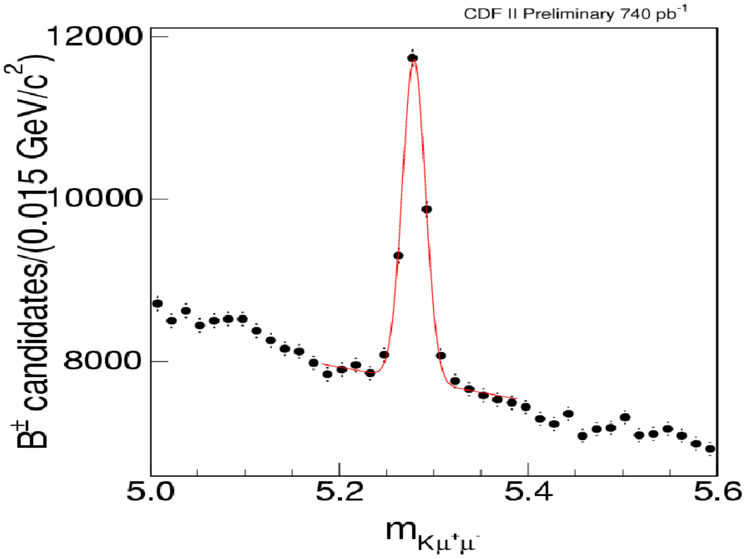
The procedure to measure the cross section is simple.

$$\sigma = \frac{N_{Data} - N_{Background}}{Acc \int Ldt}$$

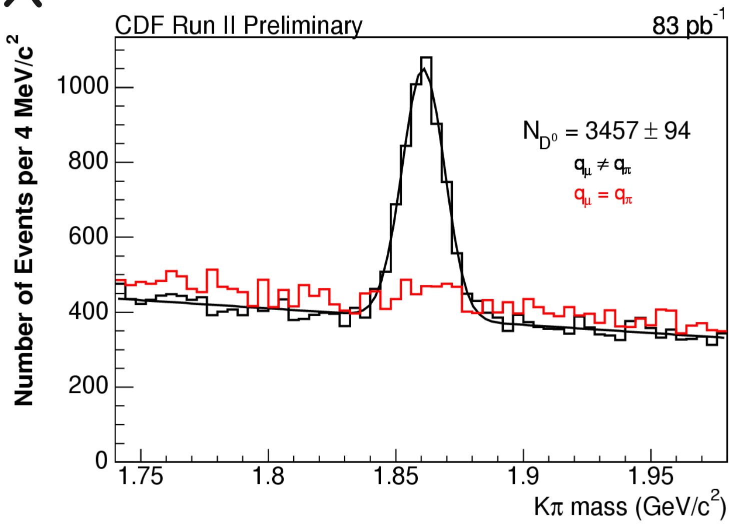
B-mesons are selected exploiting the decay channels.

The selection procedure depends on the decay channel as the background evaluation. Examples:

B →J/ψ K



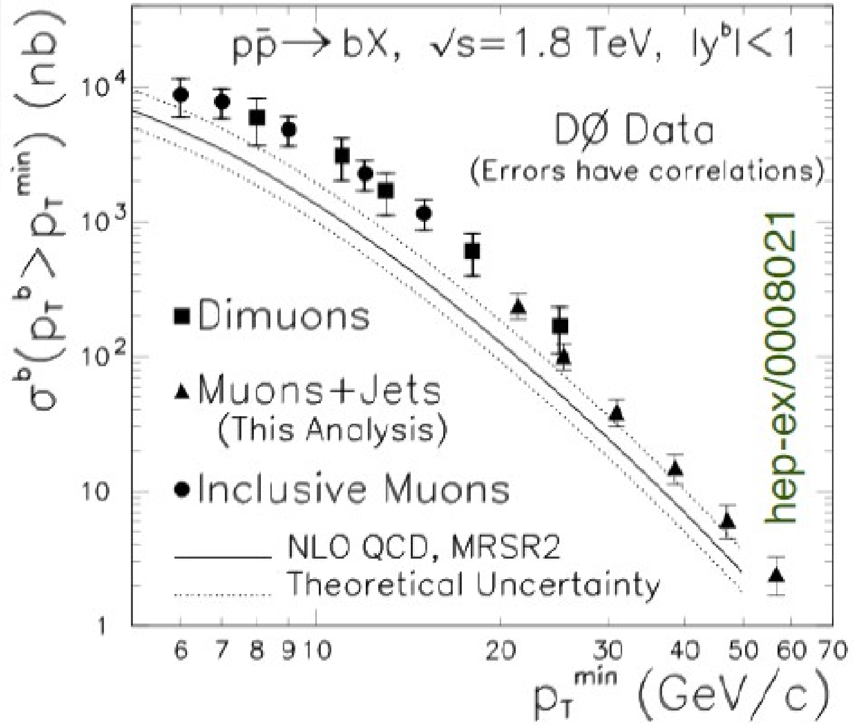
Hb → μD^{0}X

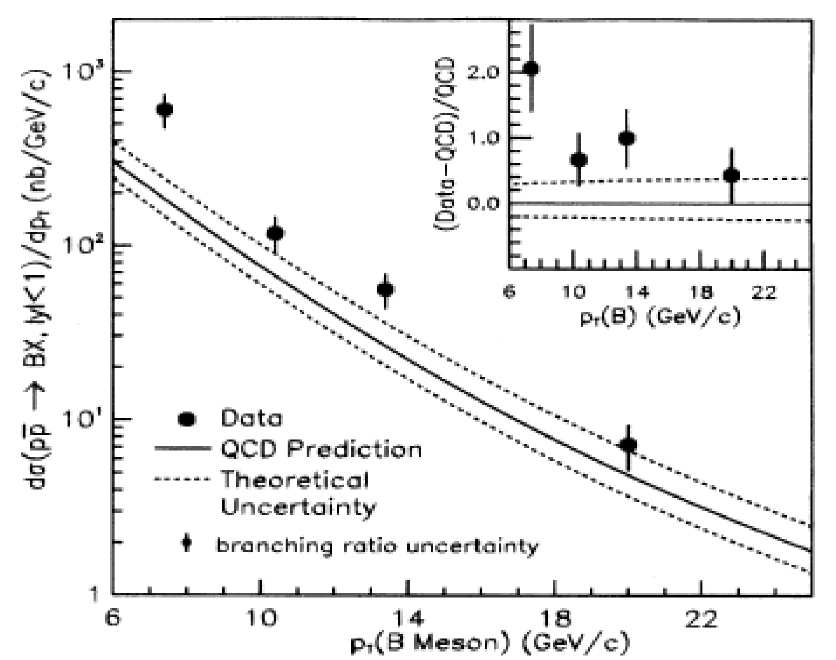


b Meson Cross Sections History

Until 2002 data and theory had a discrepancy, the measured cross section was higher of about factor 3

New Physics?





b Meson Cross Sections Measurements

A lot of work done by the experiments to improve the measurements.

A lot of work done by theoreticians to improve the theory:

M.Cacciari and P.Nason, PRL 89, 122003 (2002)

- new calculation of the cross section to include corrections (instead of NLO NLL) at high order
- new tuning of the fragmentation function D(z)
- new PDF

The purpose of this Letter is precisely to implement correctly the effect of heavy quark fragmentation in the QCD calculation. Several ingredients are necessary in order to do this: (i) A calculation with resummation of large transverse momentum logarithms at the next-to-leading level (NLL) should be used for heavy quark production [21], in order to correctly account for scaling violation in the fragmentation function. (ii) A formalism for merging the NLL resummed results with the NLO fixed order calculation (FO) should be used, in order to account properly for mass effects [22]. This calculation will be called FONLL in the following. (iii) A NLL formalism should be used to extract the nonperturbative fragmentation effects from e^+e^- data [23–29].

3

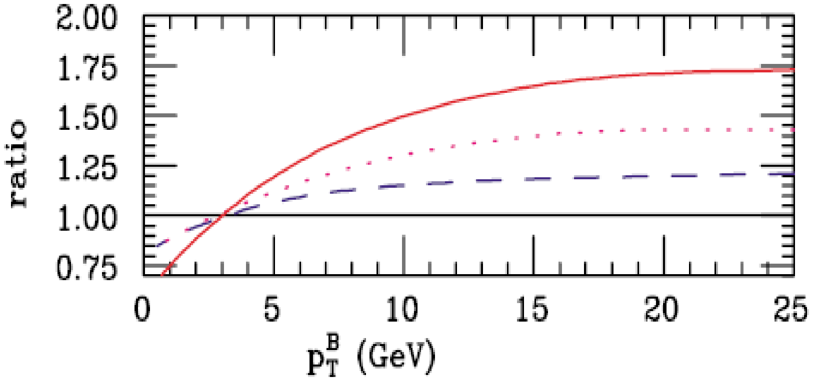
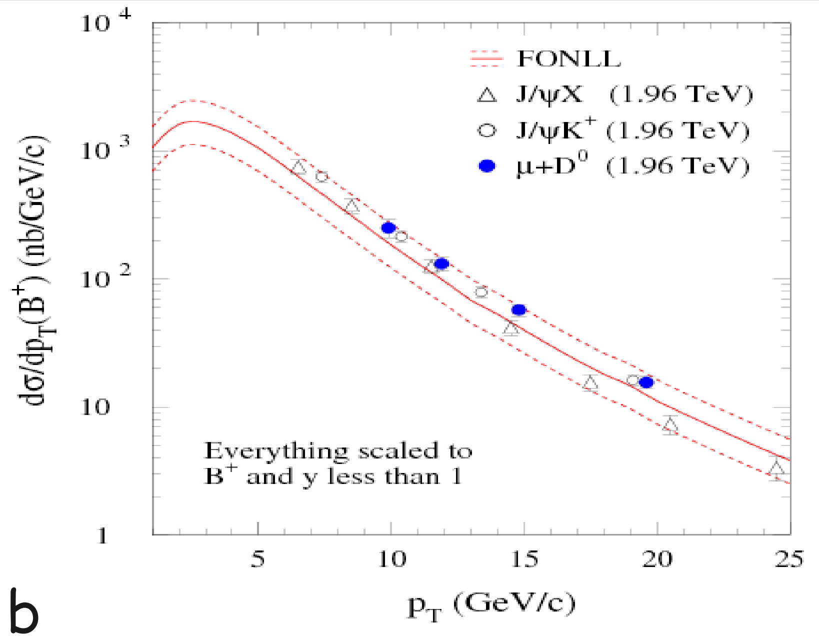


FIG. 4 (color online). The effect of the different ingredients in the calculation presented in this work, normalized to a fixed order calculation with Peterson fragmentation and $\epsilon = 0.006$. Dashed line: FO, $\epsilon = 0.002$; dotted line: FONLL, $\epsilon = 0.002$; solid line: FONLL, $N = 2$ fit.

b Meson Cross Sections Measurements

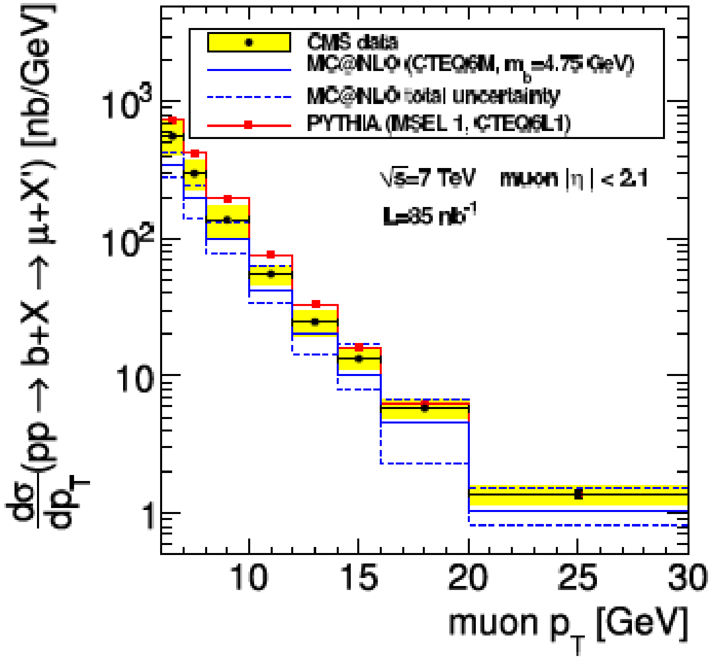
Tevatron:

- 3 separate measurements
- improved precision down to ~10%
- allow a reliable test of theory
- consistent with theory



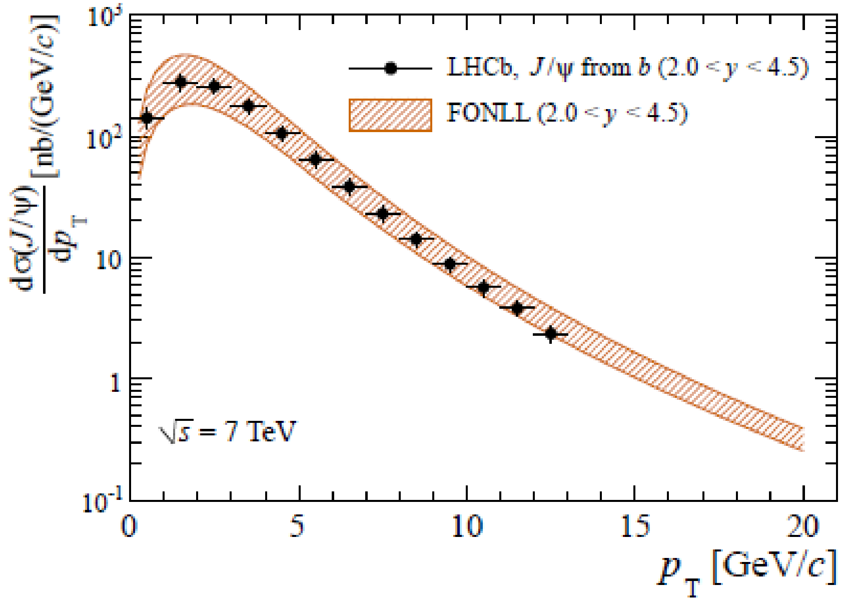
CMS:

- inclusive μ decays



LHCb:

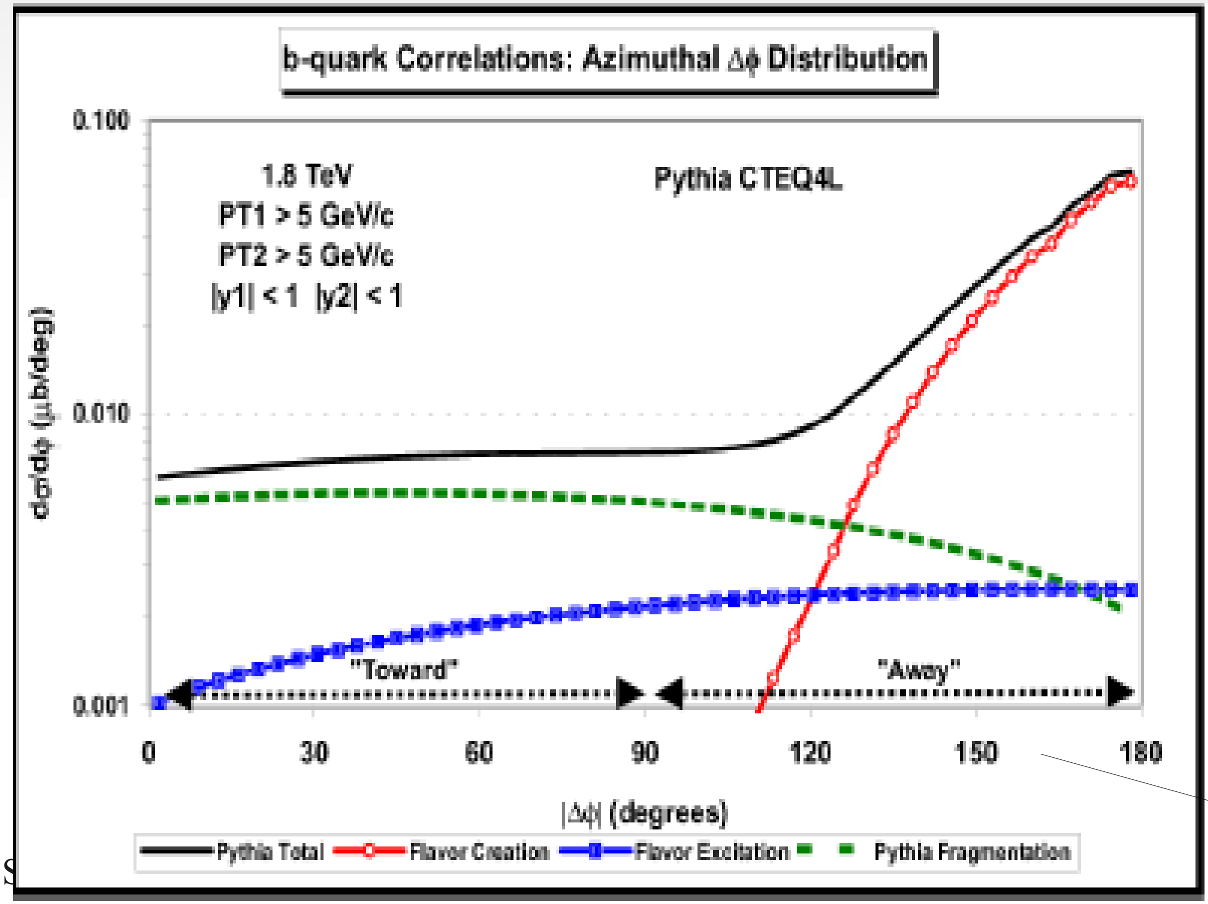
- J/ψ from b



Correlated b-$\bar{b}$ Cross Section Measurements

b-quark can be identified in jets by using inclusive selections. This allow to reconstruct events with 2 b-jets.

b-$\bar{b}$ cross section measurement allow to test high order theoretical contributions



LO produced almost back-to-back

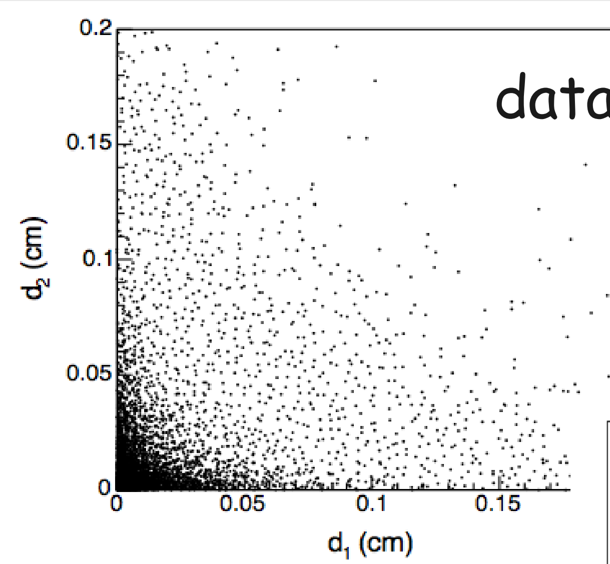
Fitting the cross section as function of the angular separation allow to determine the relative contribution of each process

ΔΦ between 2 b-jets

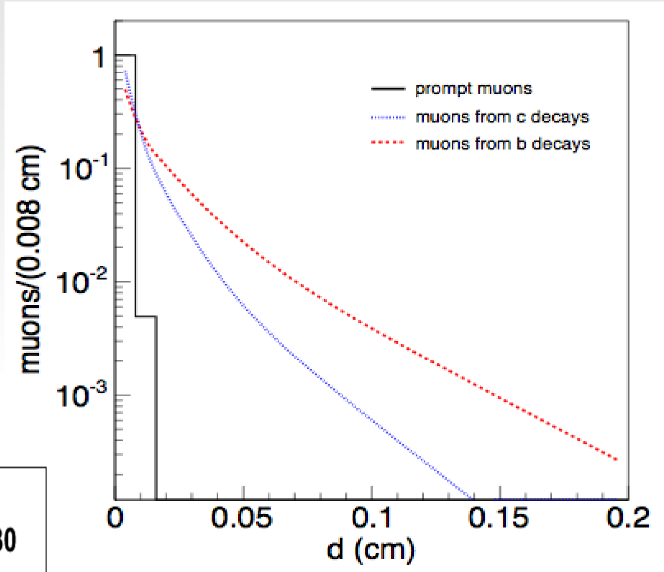
Correlated b-$\bar{b}$ Cross Section Measurements

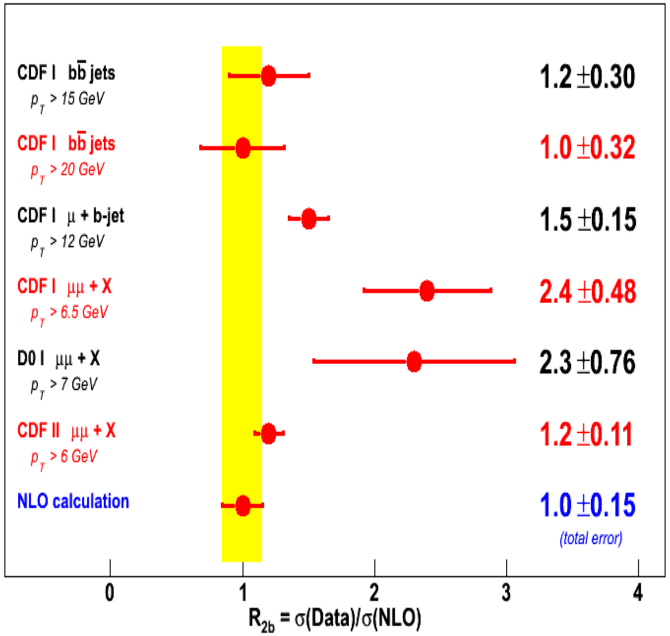
Tevatron:

2 μ are required then their impact parameter distribution is fitted to extract c and b components



Fit component



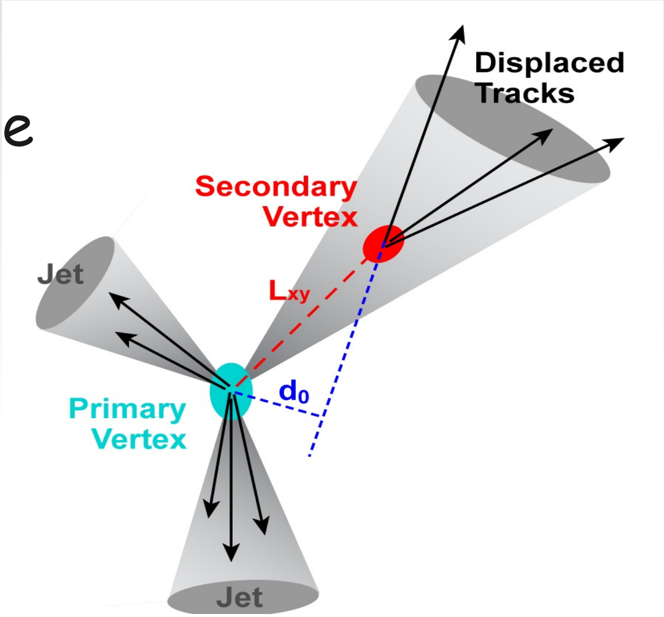


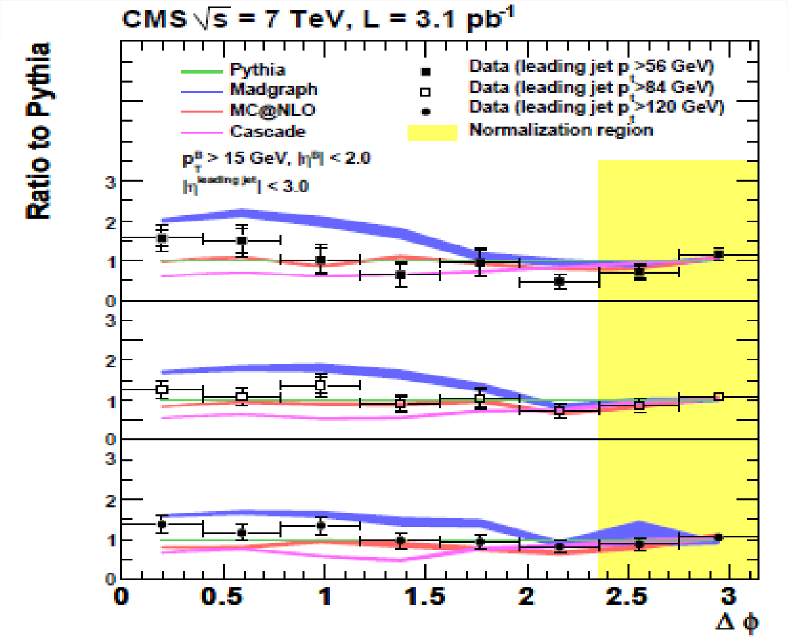
Correlated b-$\bar{b}$ Cross Section Measurements

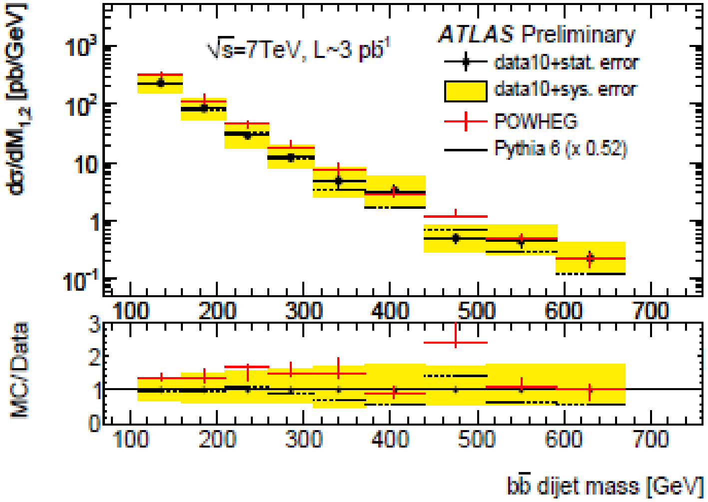
LHC:

ATLAS and CMS identify jets with b exploiting the long b lifetime

B-tagging: 2 or 3 tracks displaced from the primary vertex the decay length L_{xy} compatible with the distance traveled by the b-hadron.





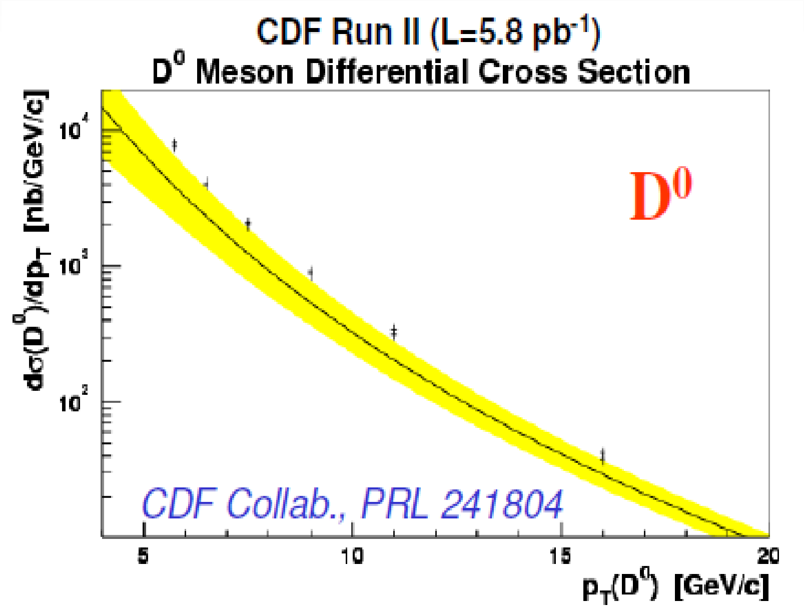


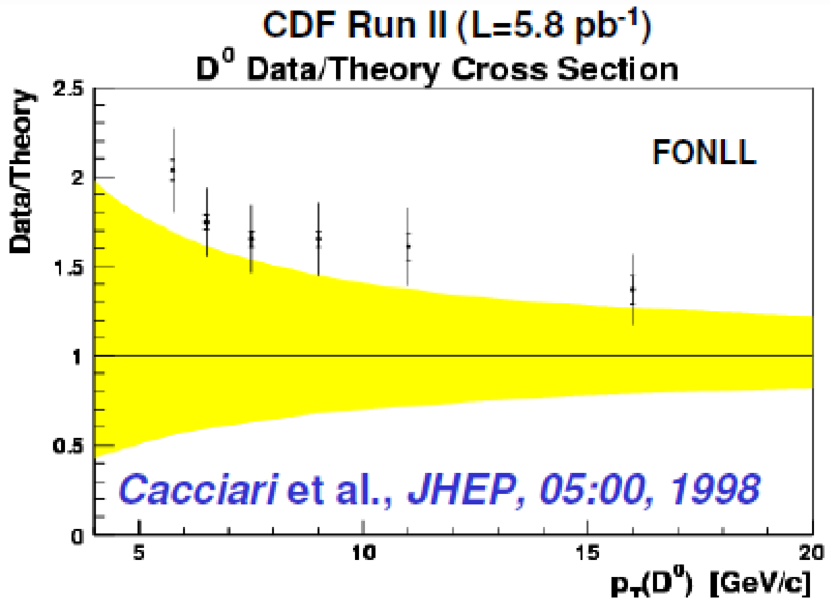
Inclusive charm cross section

Identify only one c-meson and measure the cross section with the same procedure described for b-mesons.

Theory predictions have larger uncertainties than the b-meson cross Section.

Charm meson cross section measured in several experiments, fixed target and at collider.





At Tevatron data/theory~2

Inclusive charm cross section

Also in this case a lot of work done by theoreticians

`Charm cross sections for the Tevatron Run II`
`Journal of High Energy Physics Volume 2003 JHEP09(2003)`

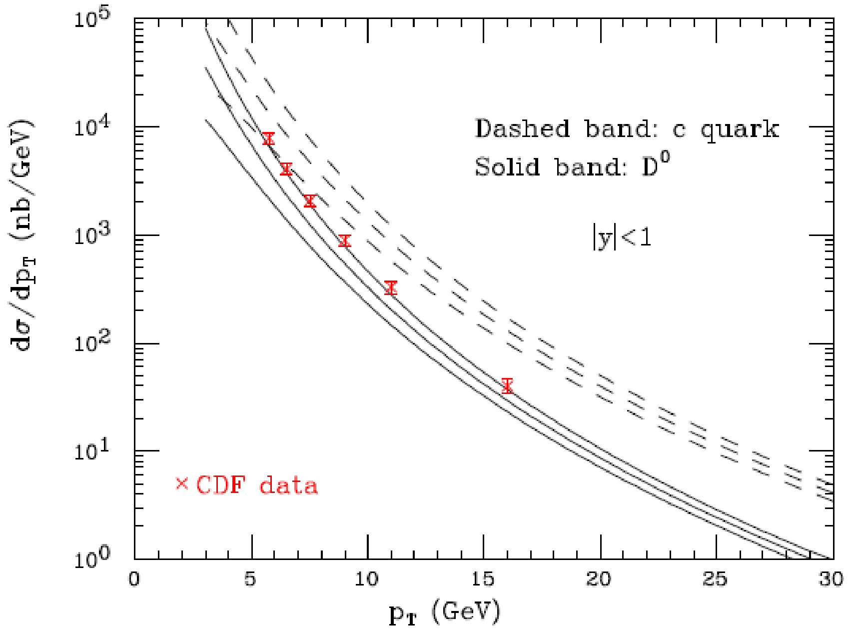
- Use the same framework of the b-meson
- Build new fragmentation functions non-perturbative using data;

Ex. $F(c \to D^0) = F_p(c \to D^0) + F(c \to D^{*+}) \otimes F(D^{*+} \to D^0) + F(c \to D^{*0}) \otimes F(D^{*+} \to D^0)$

$F(c \to D^0)$= directly produced D^0 fragmentation function

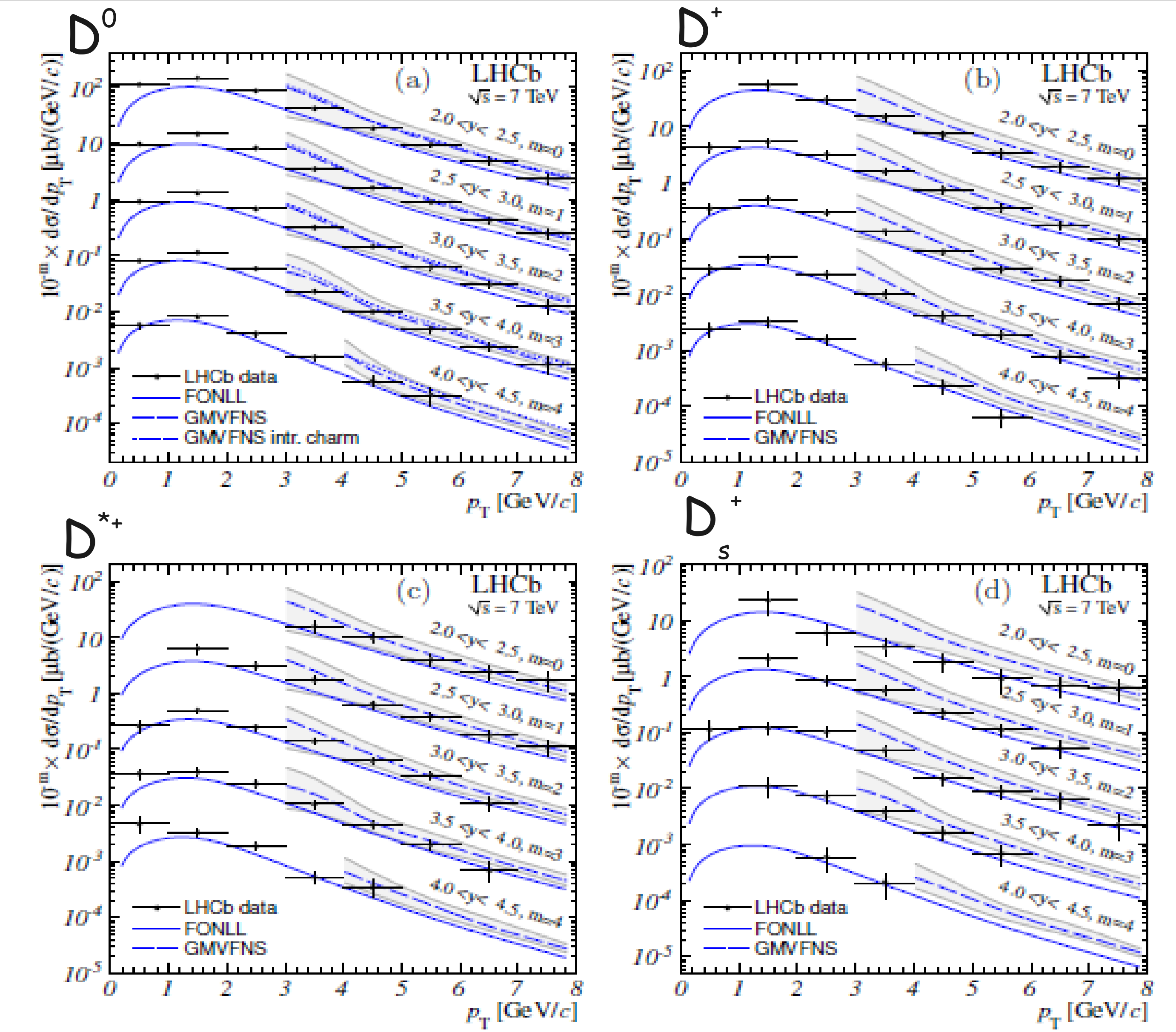
$F(c \to D^{0*})$= fragmentation function of D^{0*} that then decay to D^0

$F(c \to D^{+*})$= fragmentation function of D^{+*} that then decay to D^0

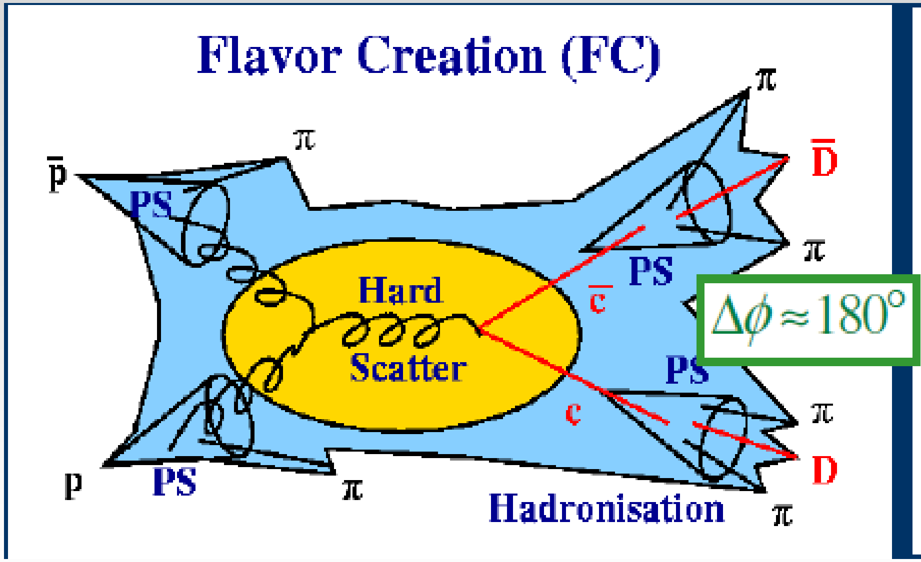


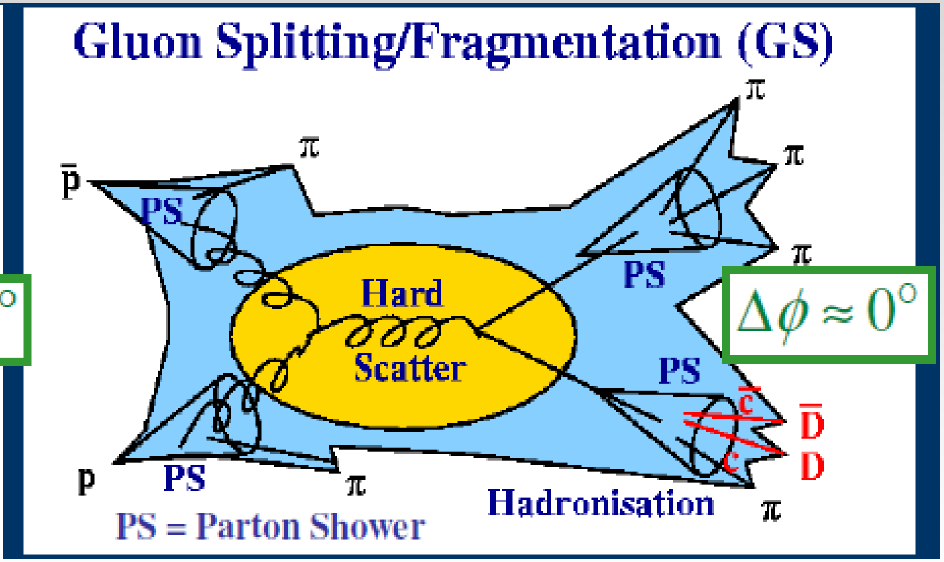
Inclusive charm cross section

Now at LHCb

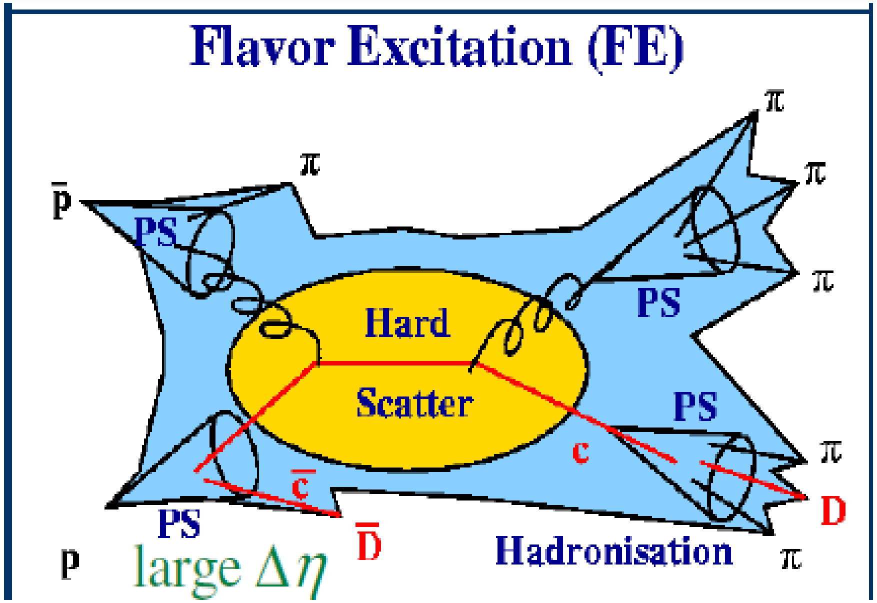


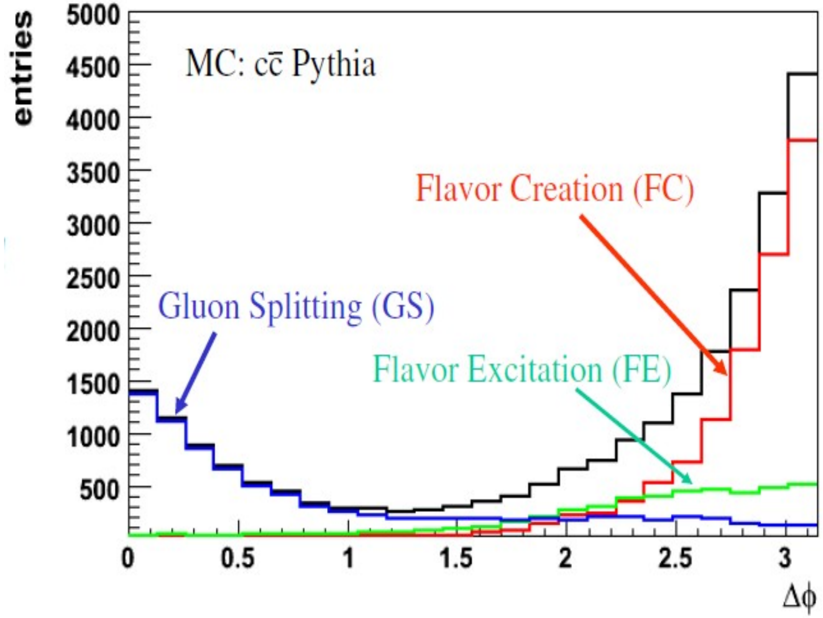
Charm Correlated Cross Sections Measurements


Flavor Creation (FC)

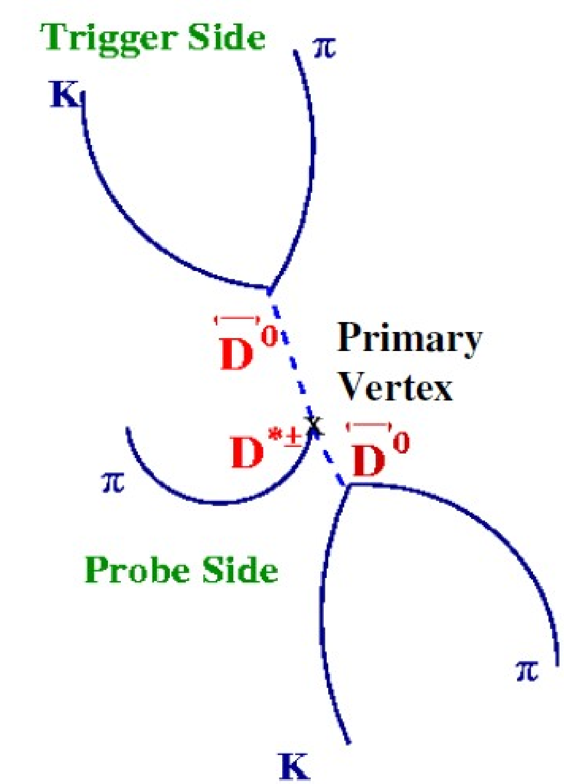

Gluon Splitting/Fragmentation (GS)

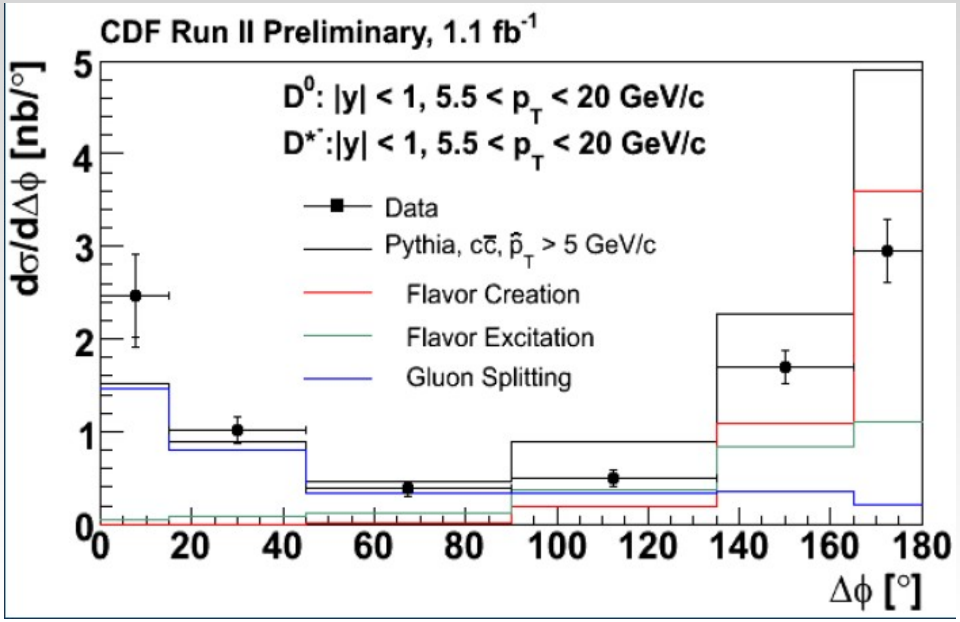
PS = Parton Shower

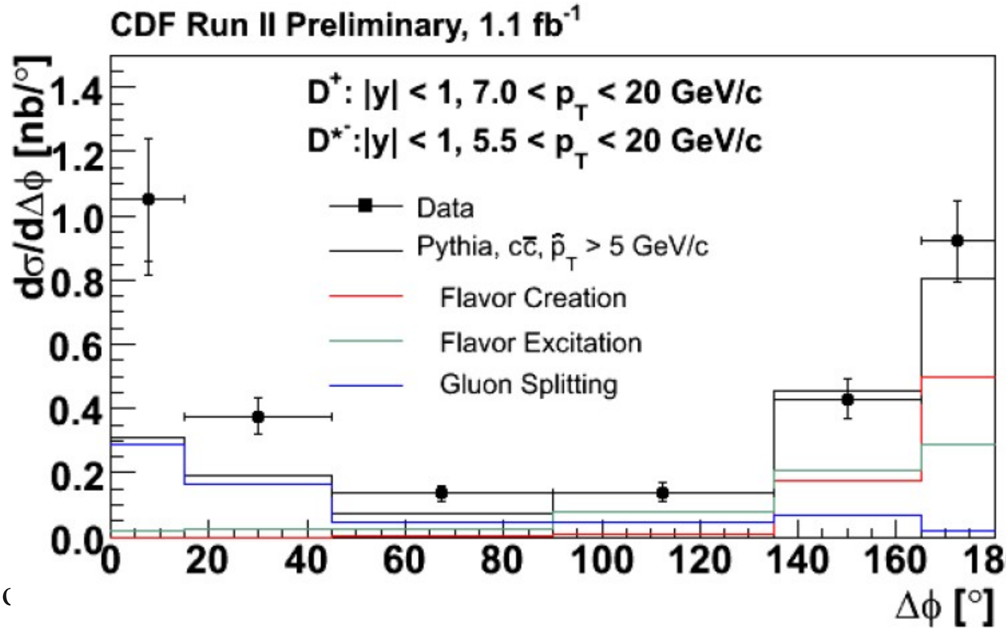

Flavor Excitation (FE)



Charm Correlated Cross Sections Measurements







- Identify the 2 D meson
- Correct for D coming from b-hadron
- Evaluated the efficiencies

Top Quark Introduction

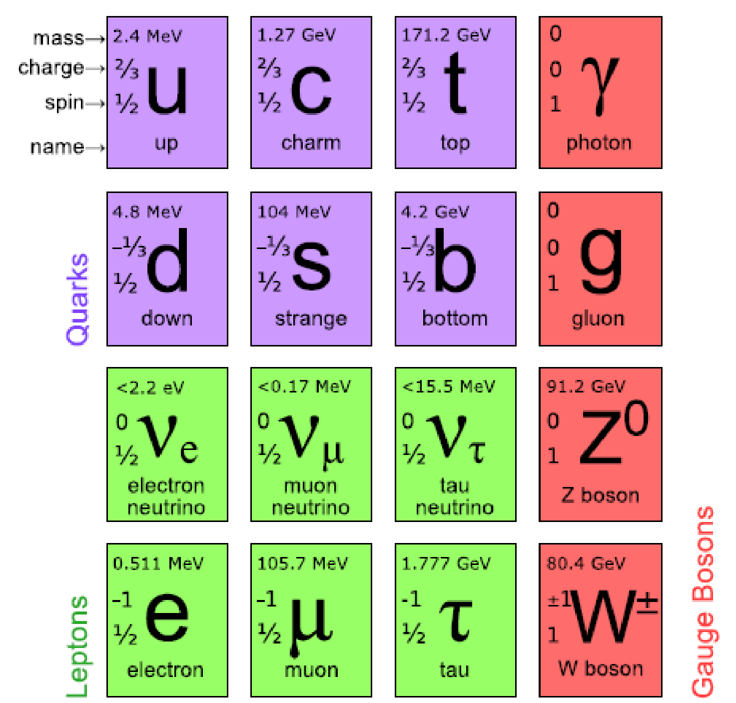
The last quark discovered. Precision SM measurements predict its existence and its mass.

In particular the asymmetry backward-forward of b-jets produced in e+e- annihilation at the Z resonance can be easily explained assuming that the b quark is in an SU(2) doublet with the top quark

Precision electroweak fits constrained the mass: 178^{+8+17}_{-8-20} GeV

The top discovery dates 1995 by the two experiments at the Tevatron Collider.

We are now in the era of precision top measurements



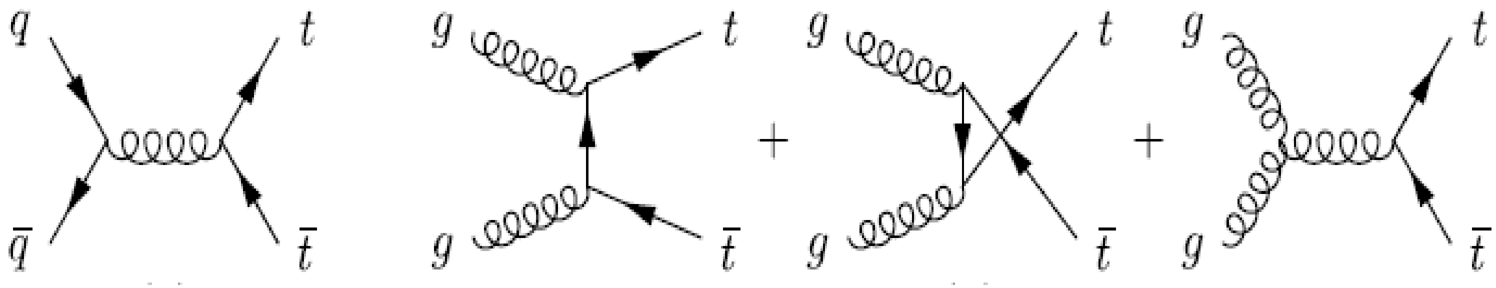
2.4 MeV ⅔ ½ u up	1.27 GeV ⅔ ½ c charm	171.2 GeV ⅔ ½ t top	0 0 1 γ photon
4.8 MeV -⅓ ½ d down	104 MeV -⅓ ½ s strange	4.2 GeV -⅓ ½ b bottom	0 0 1 g gluon
<2.2 eV 0 ½ ν_e electron neutrino	<0.17 MeV 0 ½ ν_μ muon neutrino	<15.5 MeV 0 ½ ν_τ tau neutrino	91.2 GeV 0 1 Z^0 Z boson
0.511 MeV -1 ½ e electron	105.7 MeV -1 ½ μ muon	1.777 GeV -1 ½ τ tau	80.4 GeV ±1 1 $W^\pm$ W boson

Top Quark Cross Sections

$$\sigma(pp \to t\bar{t} + X) = \sum_{i,j} \int dx_i dx_j \times F_i(x_i,\mu) F_i(x_j,\mu) \hat{\sigma}_{ij}(x_i,x_j,m_{top}^2,\mu^2)$$

$m_{top}/2 < \mu < 2m_{top}$ since the mass is so large the calculation can be performed with the perturbative QCD

At LO the diagrams that contribute are



LHC: 80% gluon fusion 20% $q\bar{q}$
Tevatron: 85% $q\bar{q}$ 15% gluon fusion

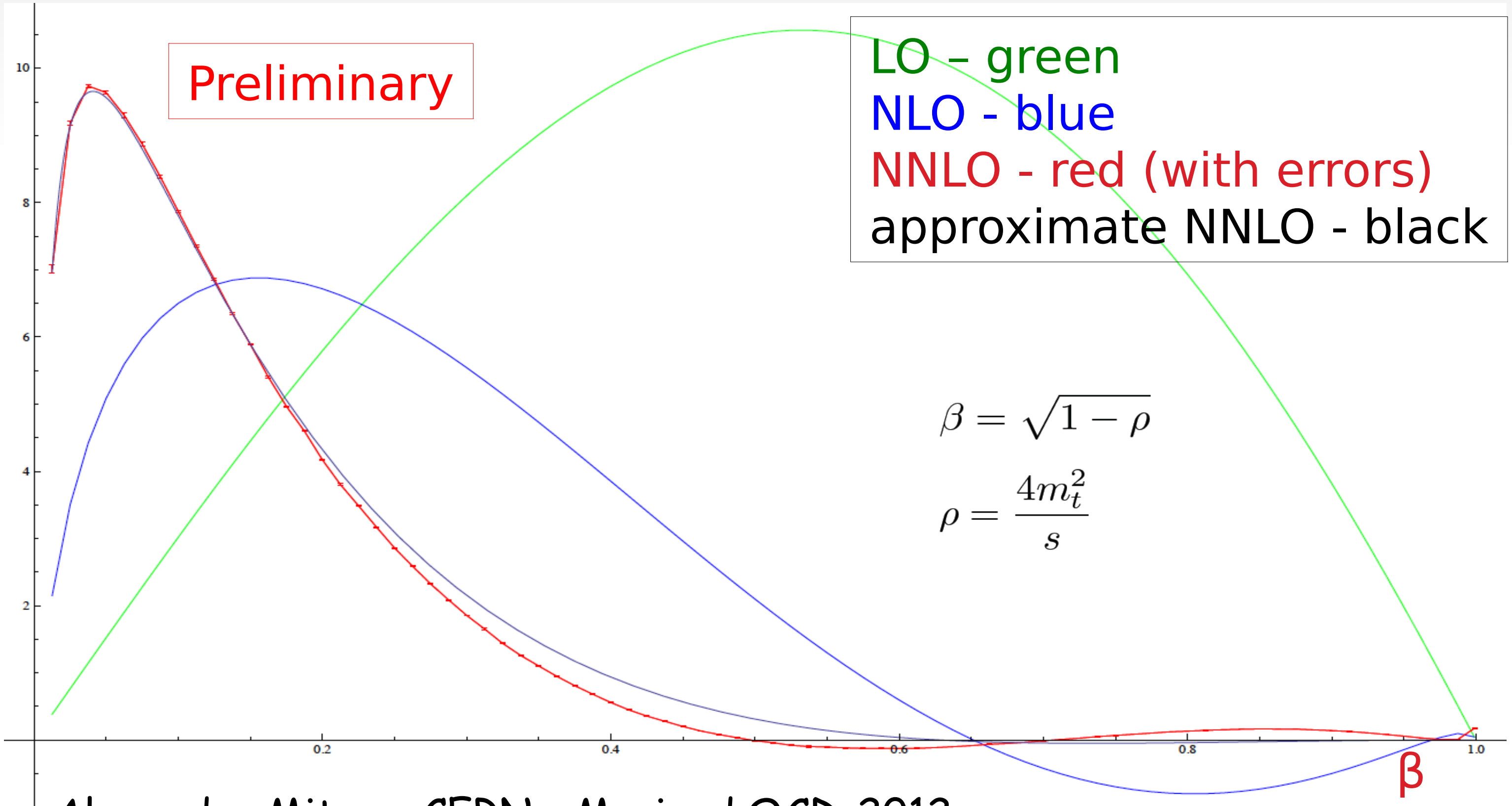
NLO calculations available.

Top Quark Cross Sections high order

NLO calculations are important: ~50%

Since not everything is in agreement with the theoretical expectations theoreticians are calculating also the NNLO corrections

$$\hat{\sigma}_{q\bar{q}\to t\bar{t}}(\beta) = \frac{\alpha_S^4(m_t)}{m_t^2}\left\{\mathrm{LO} + \mathrm{NLO} + \mathrm{NNLO}\right\}$$



$$\beta = \sqrt{1-\rho}$$

$$\rho = \frac{4m_t^2}{s}$$

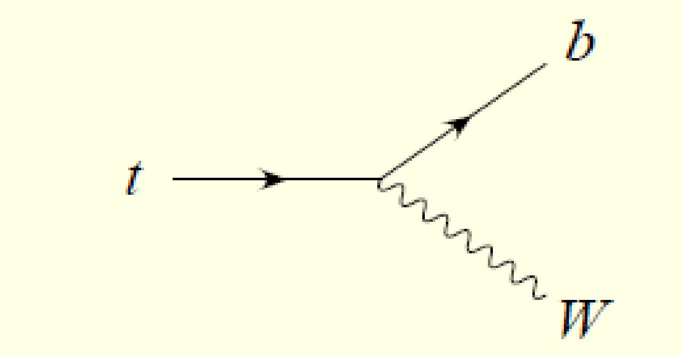
Alexander Mitov – CERN - Moriond QCD-2012

Top Quark Decay

Quark top decay before it can form a bound state

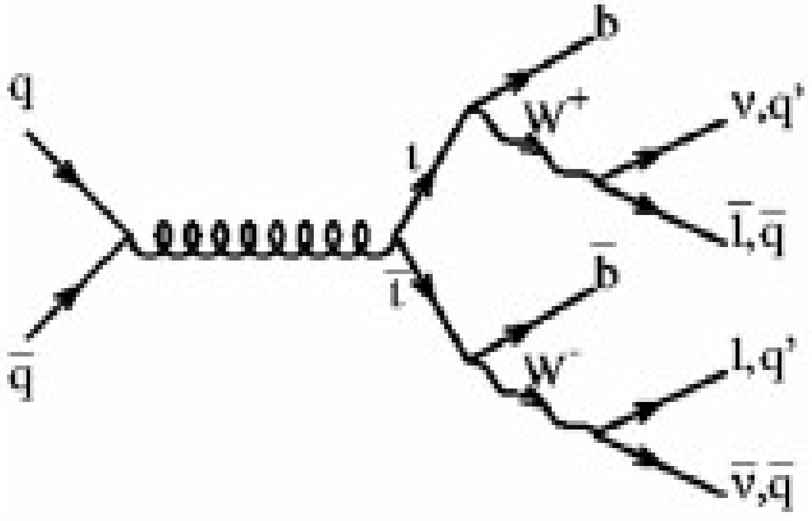
$\tau_t \simeq 10^{-25}$ sec compare to $\tau_{\mathrm{QCD}} \simeq 10^{-24}$ sec

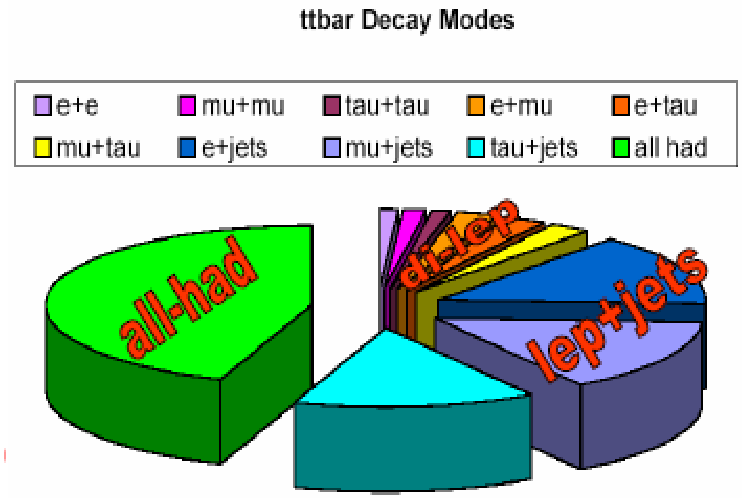
It decays predominantly



$$t \to bW^+ \begin{cases} W^+ \to l^+\nu_l \\ W^+ \to q\bar{q}' \end{cases}$$

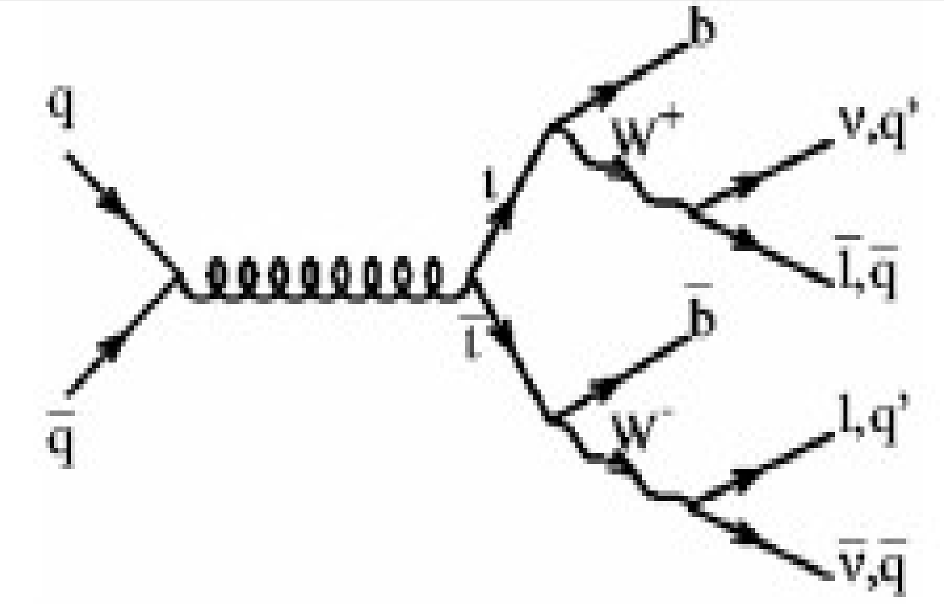
The event is



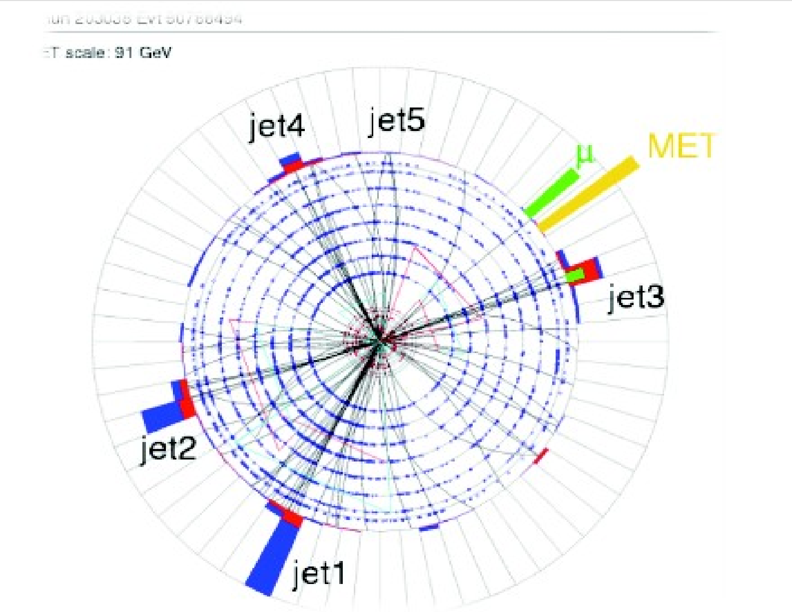


Top Quark Reconstruction

Theory



Detector



Events classified depending on the W decay:

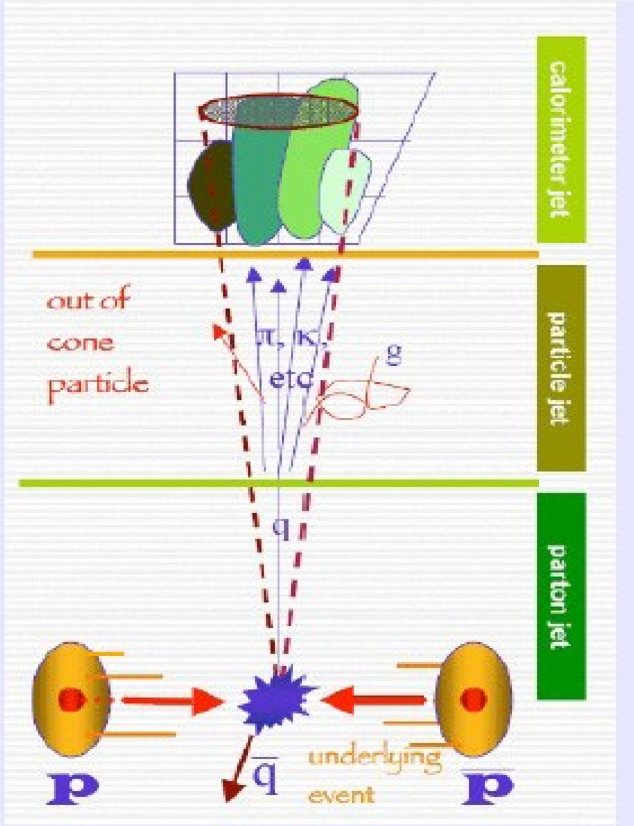
- Di-lepton: low yield, low background, well defined leptonic signature, neutrinos → MET
- Lepton+jets: higher yield, moderate background, lepton signature + MET + jets
- All hadronic: highest yield, huge background, only jets

Top Quark Events Reconstruction: Common tools

Final states always with jets and b-quark in jets.

1. Reconstruct jets
2. Use b-tag algorithm to determine if the jet is originated by a b-quark

1.



Use cone algorithm

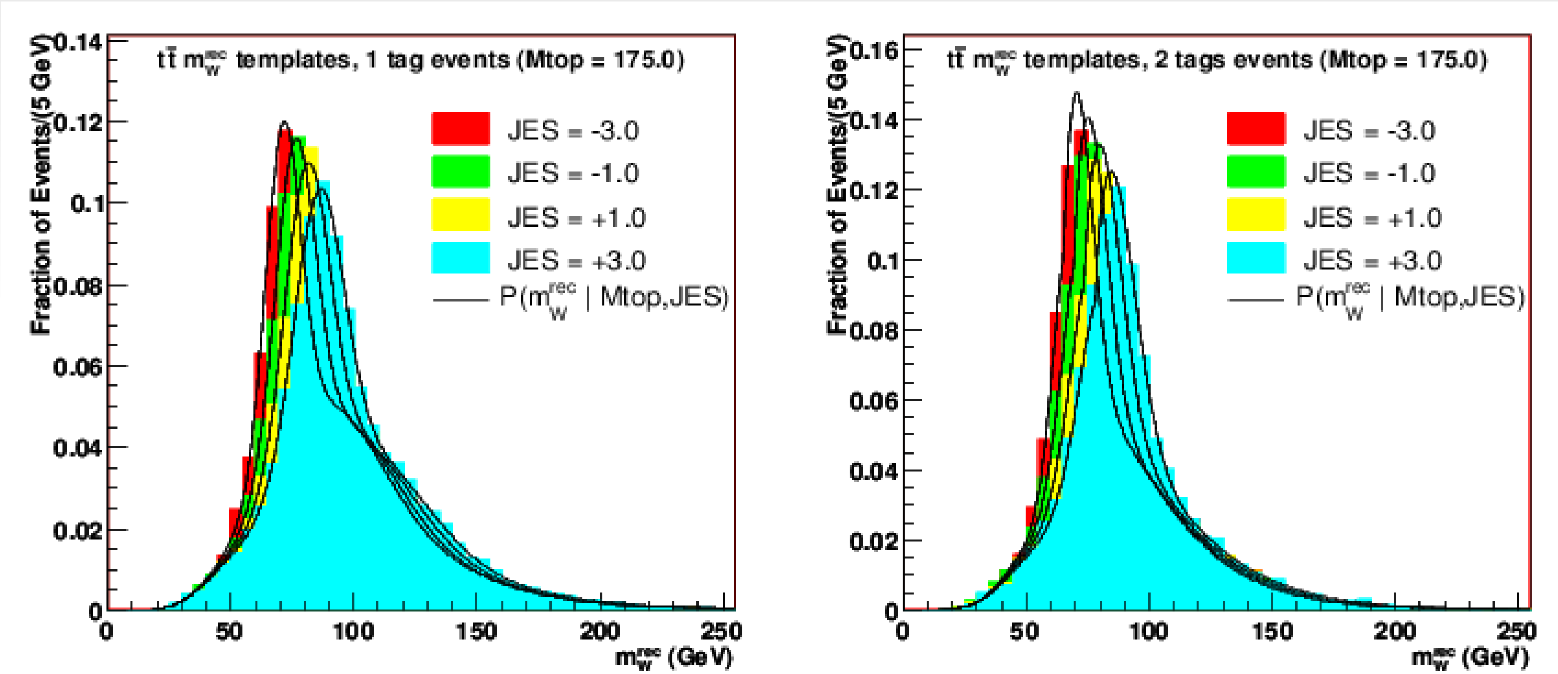
Jet Energy Scale (JES) is one of the major source of uncertainty (see discussion on jet reconstruction)

Top analysis now use a new method to determine the energy scale: the "in situ" calibration.

Common Tools: “In situ” Energy Calibration

In the decay channels where both Ws decay in hadrons it is possible to leave the JES as free parameter and fit the W mass.

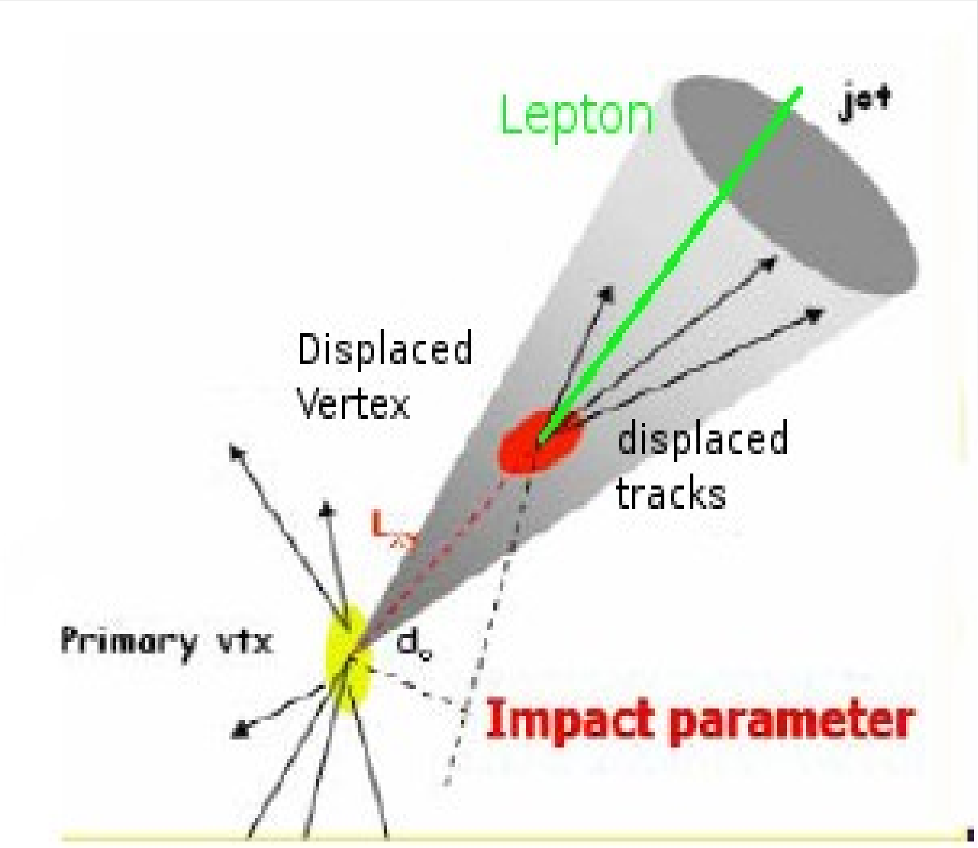
Templates with different JES are produced and the W mass is fitted



$$\chi^2 = \frac{\left(M_{jj}^{a} - M_W^{rec}\right)^2}{\Gamma_W} + \frac{\left(M_{jj}^{b} - M_W^{rec}\right)^2}{\Gamma_W}$$

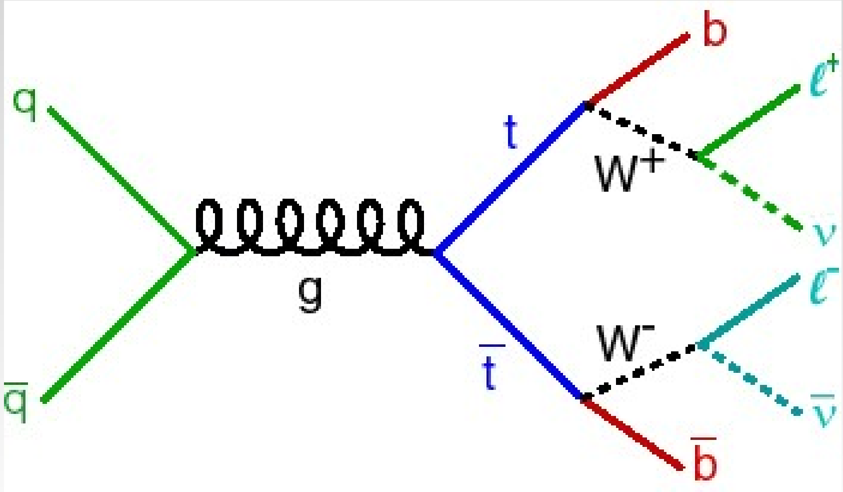
Top Quark Reconstruction:Common tools

2. Use b-tag algorithm to determine if the jet is originated by a b-quark



- Select tracks with high impact parameter respect to primary vertex
- Request at least 2 tracks
- Fit the tracks to identify a secondary vertex
- Cut on decay lenght L_{xy} to be compatible with the distance traveled by a b-hadron

Top Quark Decay Selections

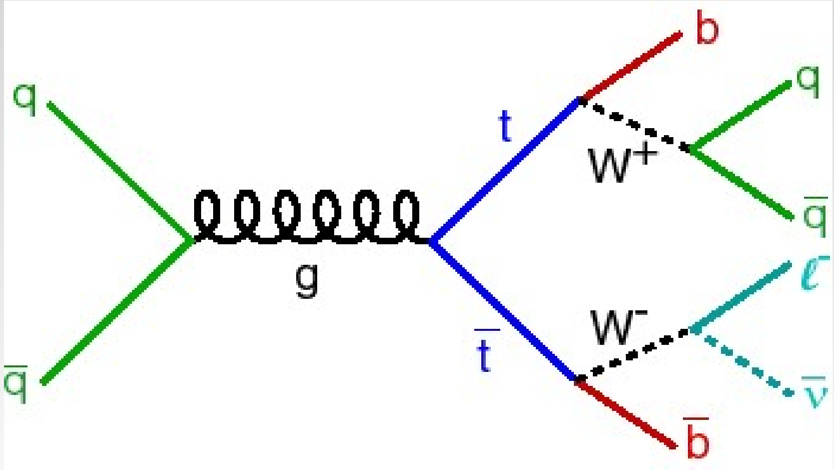


Requirements:

- two high P_T opposite charge isolated leptons
- at least 2 high E_T jets
- at least one vertex b-tag
- Significant MET

Major Backgrounds

- Process with 2 leptons in the final state: Drell-Yan Z/γ^*, WW,WZ,ZZ
- QCD: fake leptons



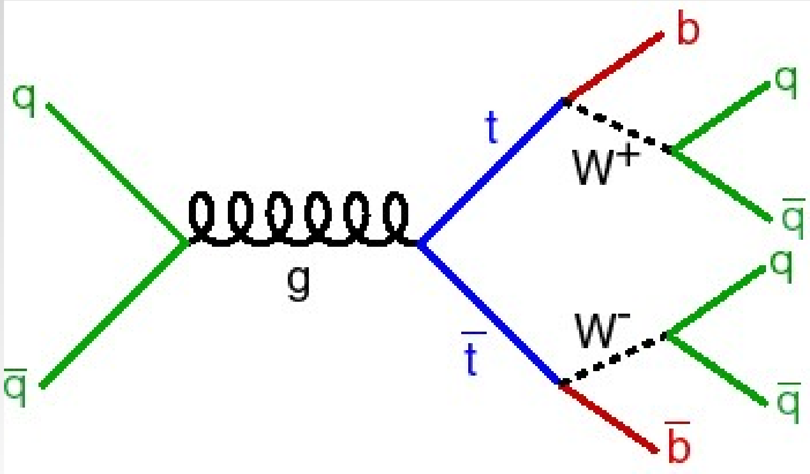
Requirements:

- one high P_T isolated leptons
- at least 4 high E_T jets
- at least one b-tag
- Significant MET

Major Background

- Process with 1 lepton + jets in the final state: W+jets
- Other contributions from non-W

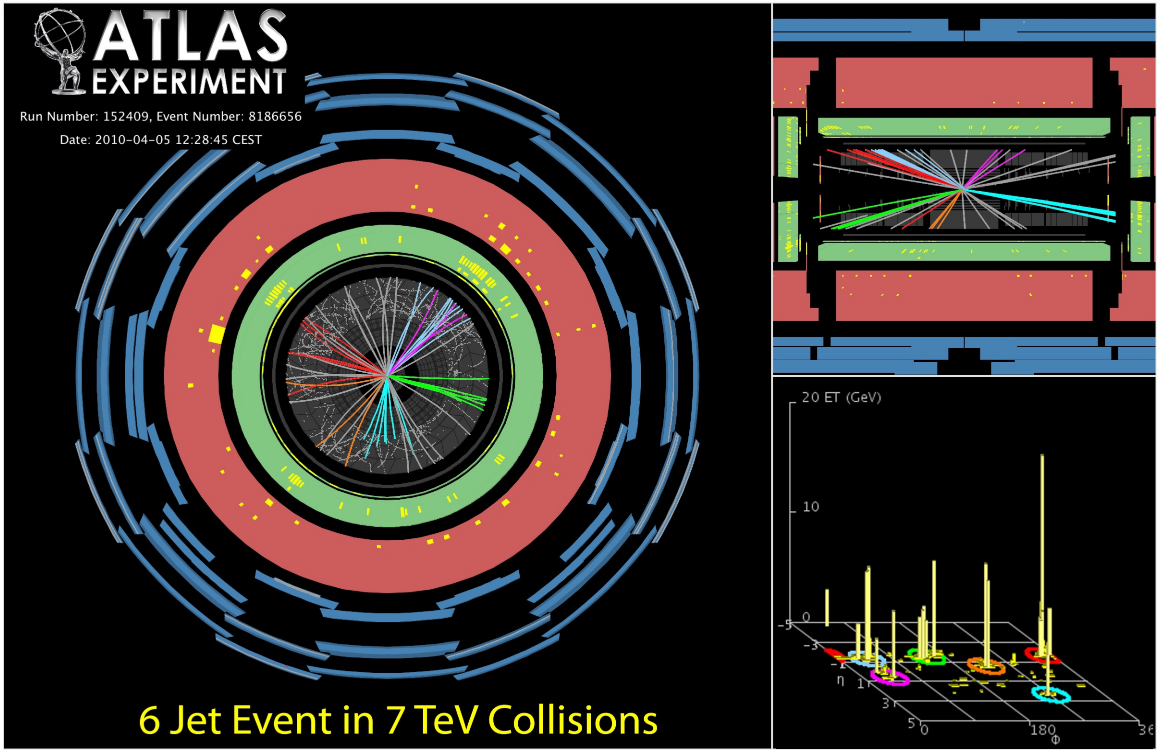
Top Quark Decay Selections



Requirements:

- at least 6 high E_T jets
- at least one b-tag
- Small MET
- No leptons

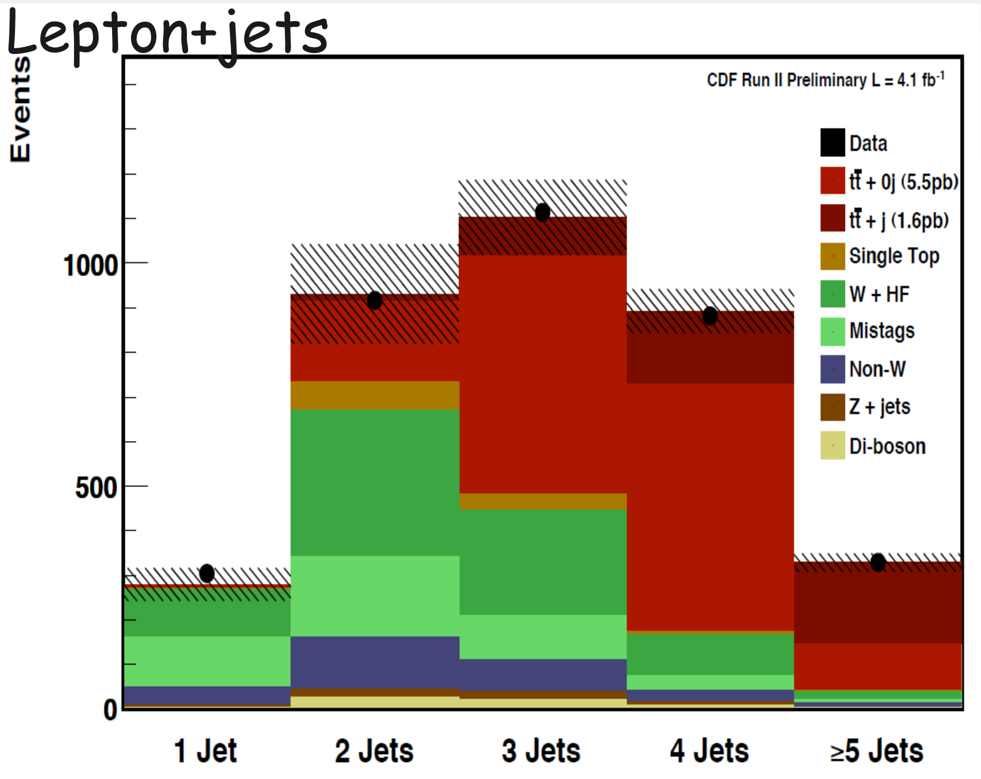
Dominant Background: QCD multi-jets

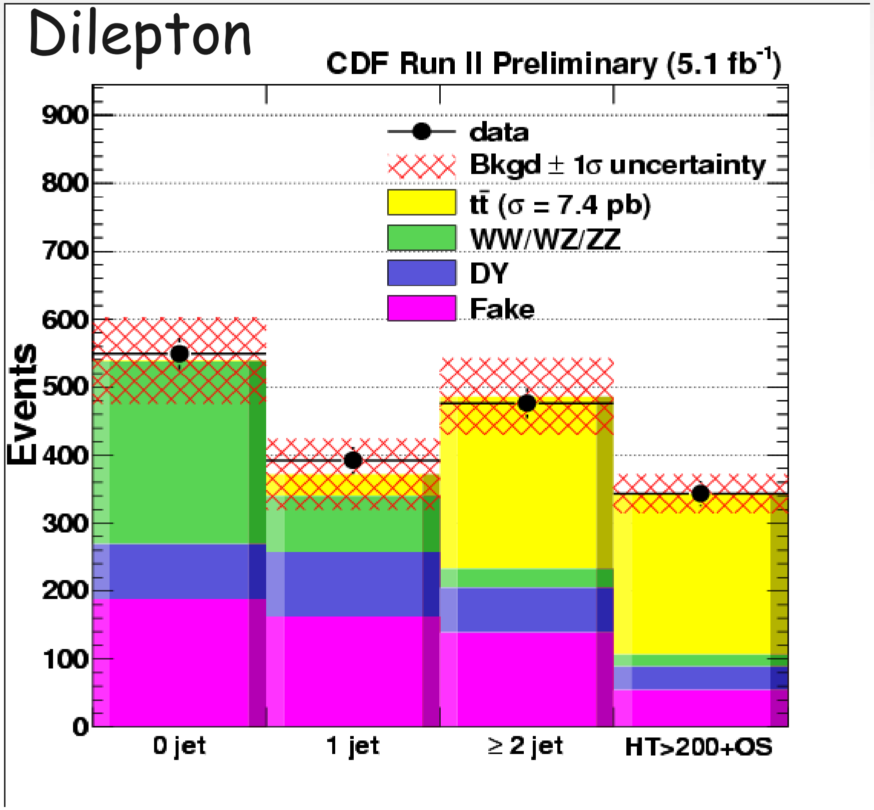


6 Jet Event in 7 TeV Collisions

Top Quark Event count

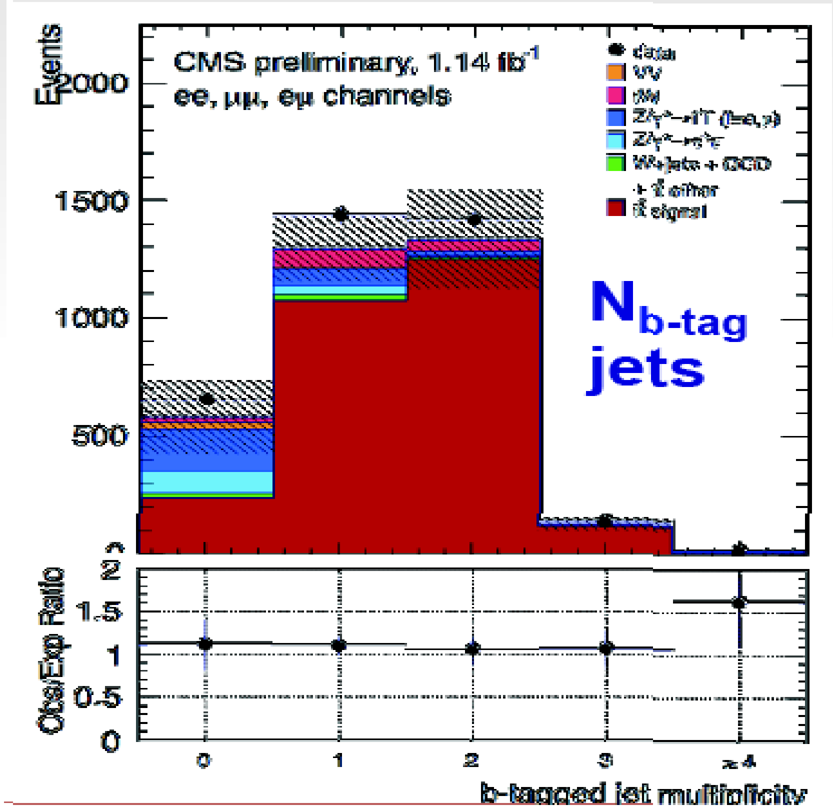
In order to count the number of top-anti-top event candidates the number of events is plotted versus the n umber of jets per event. In each bin the contribution of signal and background is different.

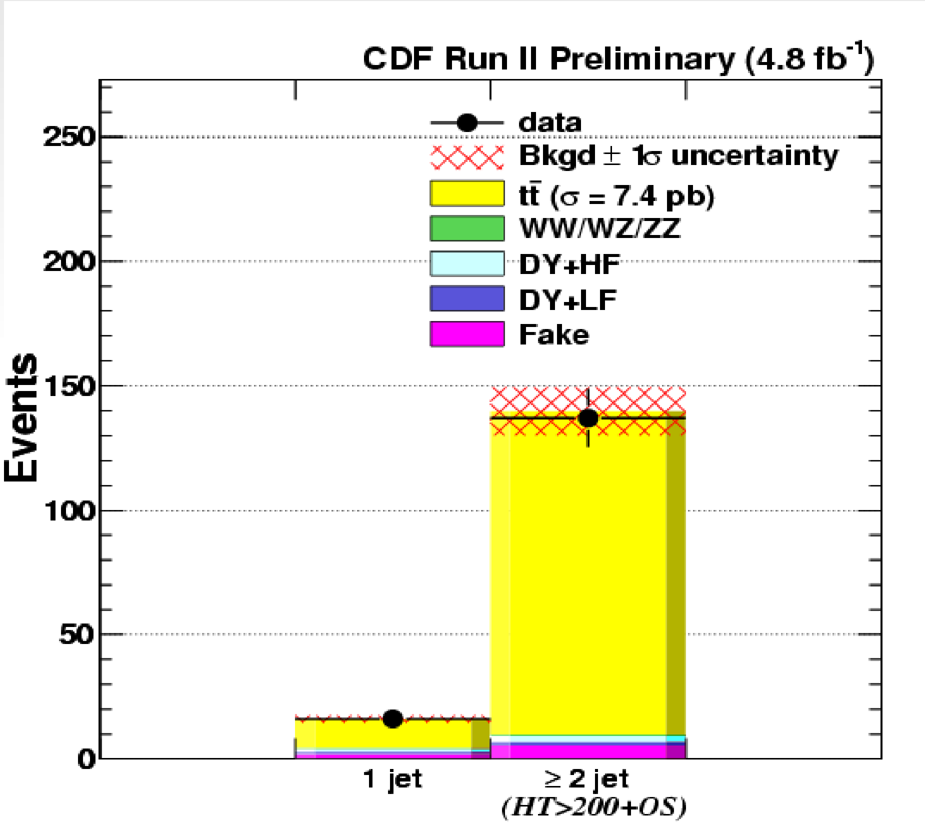

Lepton+jets


Dilepton

Top Quark Event count - 2

In order to increase the purity of the sample the number of b-tagged jets are counted or at least 2 b-jets are required.

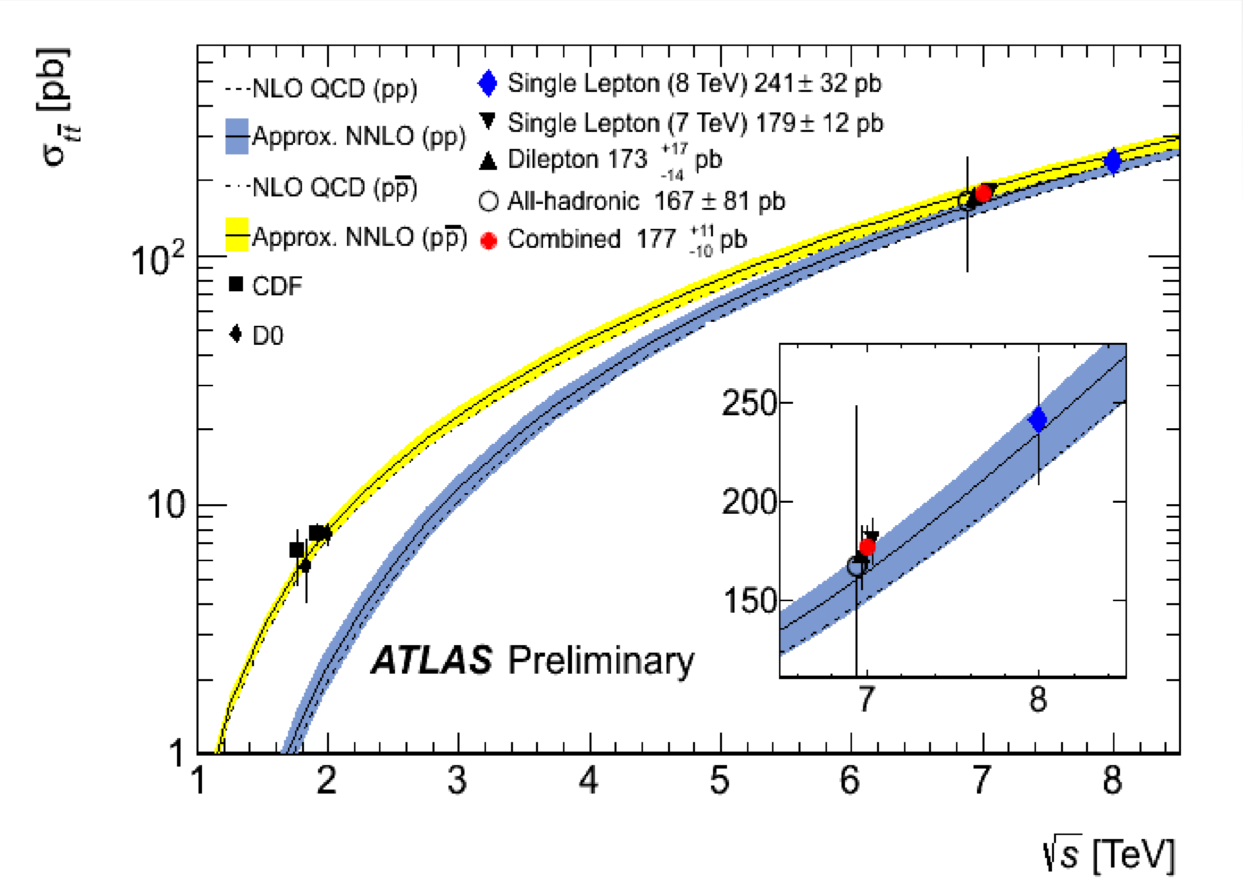




Top Quark Cross Section

$$\sigma_{tt} = \frac{N_{Data} - N_{Background}}{Acc\int Ldt}$$

Inserting the number of signal and background events in the formula and knowing luminosity and efficiency on signal we have the cross section

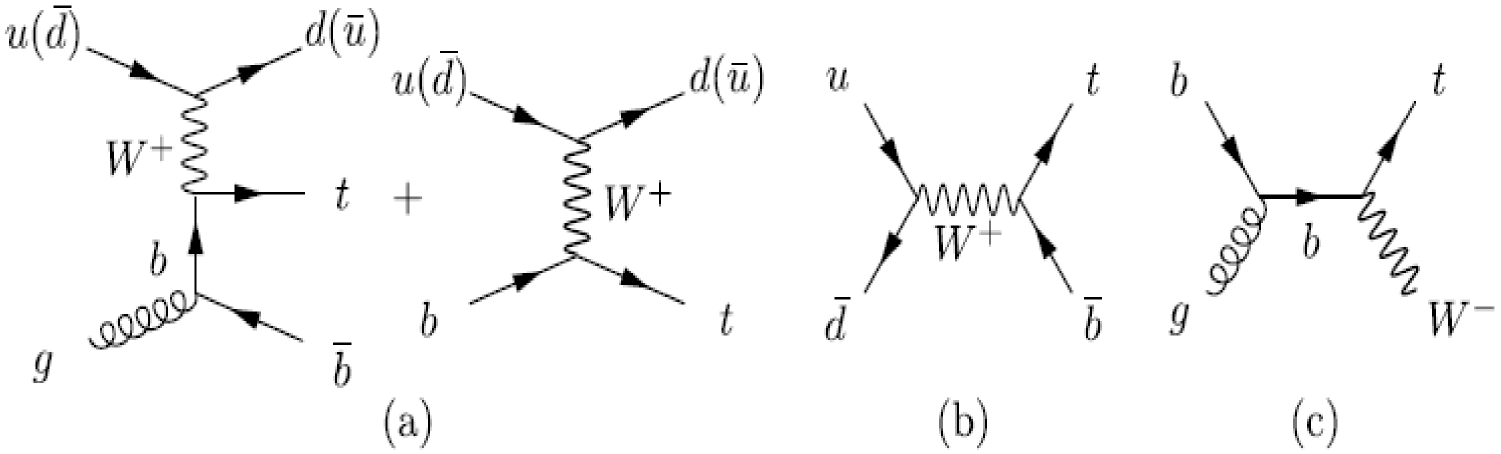


Good agreement with the expectations

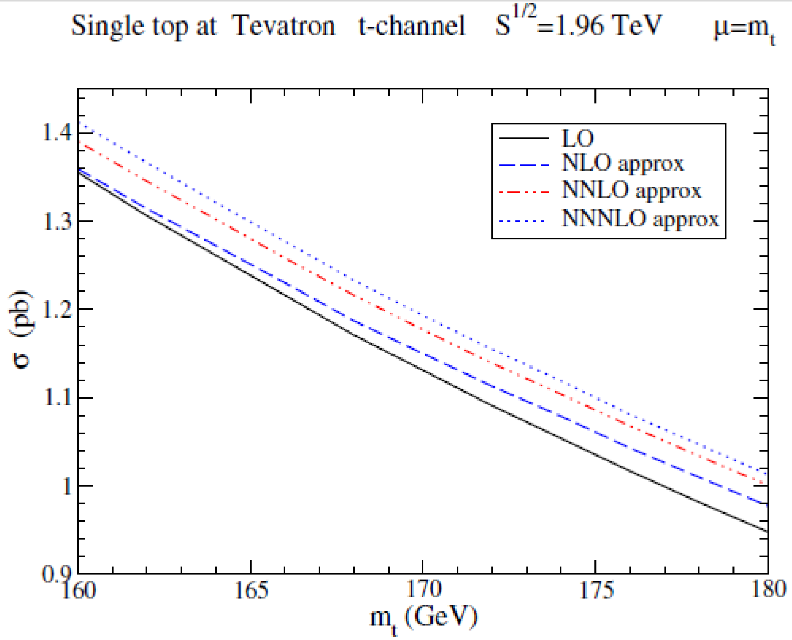
Single Top Quark

Top can be produced also via electroweak interaction involving a vertex Wtb. There are three different production models depending on the Q^2 of the W:

1. t-channel: a virtual W-boson interact with b-quark (sea quark) (a)
2. s-channel: a virtual W boson $q^2 > (m_{top} + m_b)^2$ is produced by the fusion of 2 quark of SU(2) isospin doublet (b)
3. W-associated production: top quark is produced with a real W-boson starting from a sea b-quark and gluon (c)



Single Top Quark Expected Cross Section

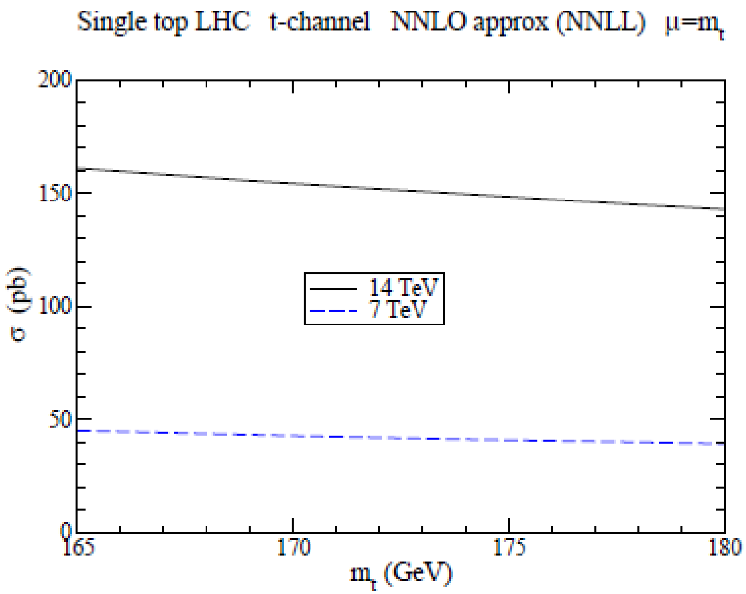


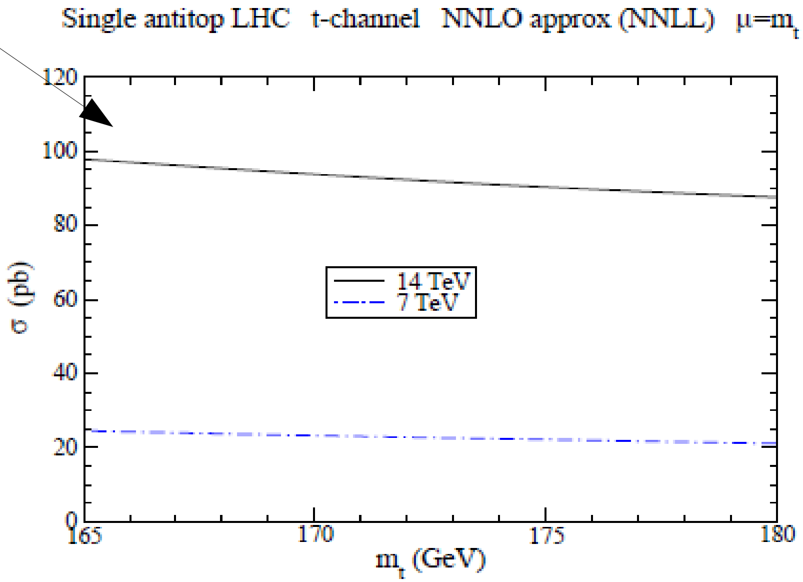
Single top cross section t-channel

Dominate at Tevatron and LHC

PRD74,114012,(2006)

LHC

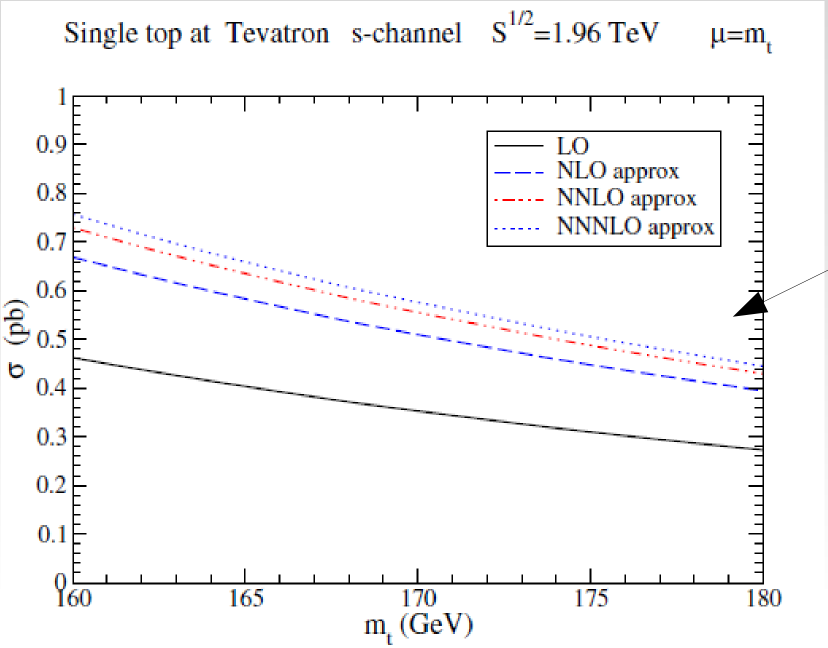




N. Kidonakis, DIS 2011, Newport News, Virginia, April 2011

Single Top Quark Expected Cross Section

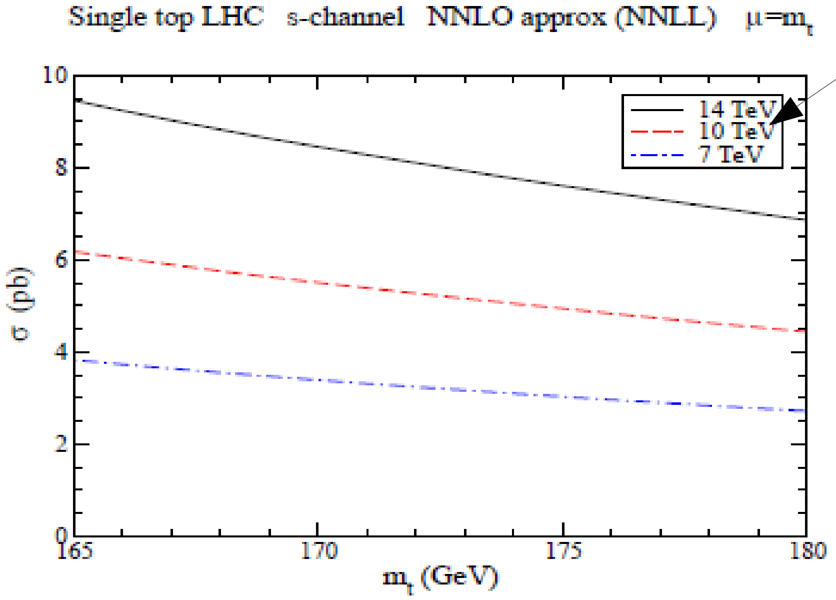
Single top cross section s-channel

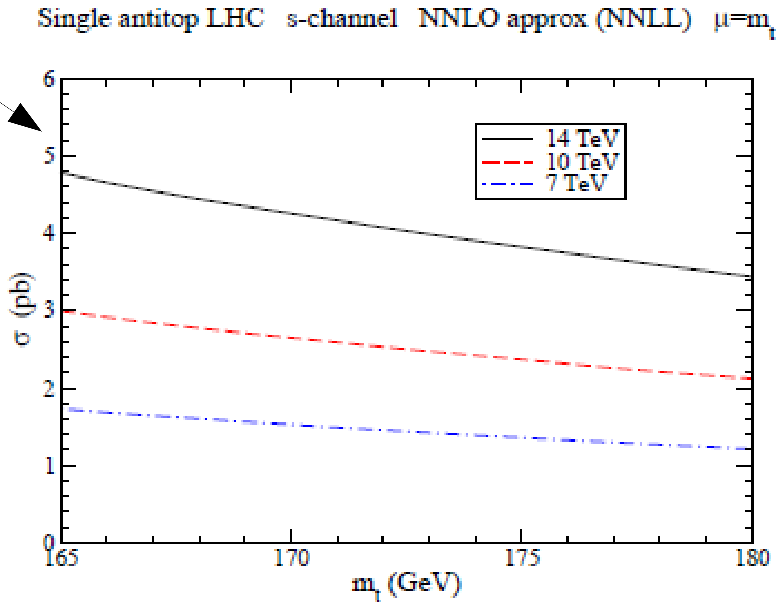


Tevatron

PRD74,114012,(2006)

LHC

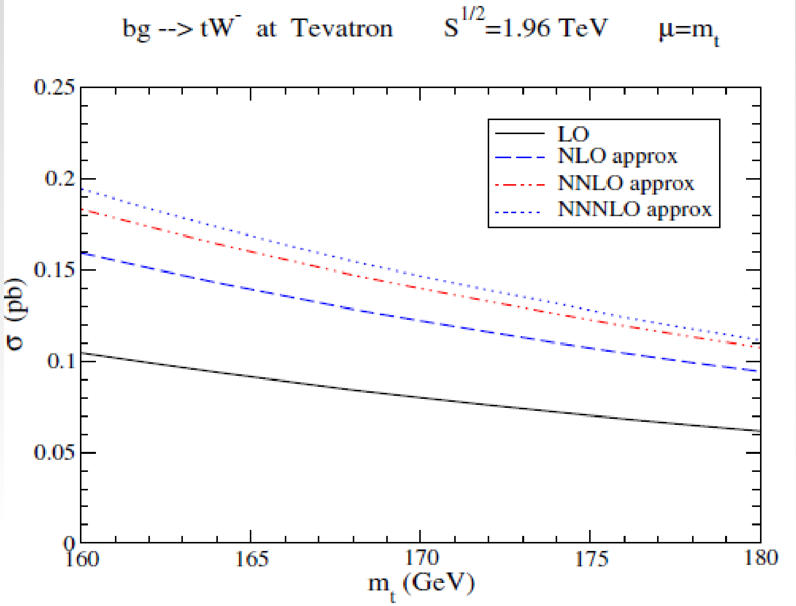


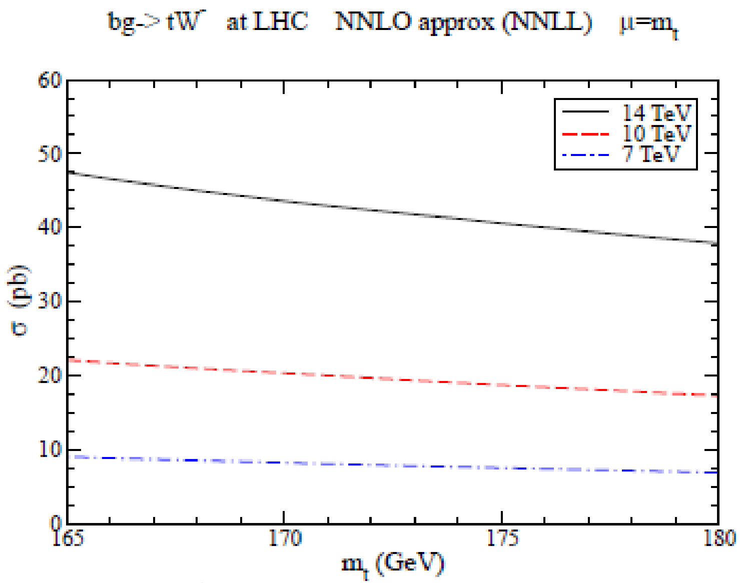


arXiv:1005.3330

Single Top Quark Expected Cross Section

PRD74,114012,(2006)





arXiv:1005.3330

Single top cross section:

Wt associated production

Not enough sensitivity at Tevatron

Single Top Quark Expected Cross Section